

# Coulomb Sturmiian basis functions in electronic structure theory

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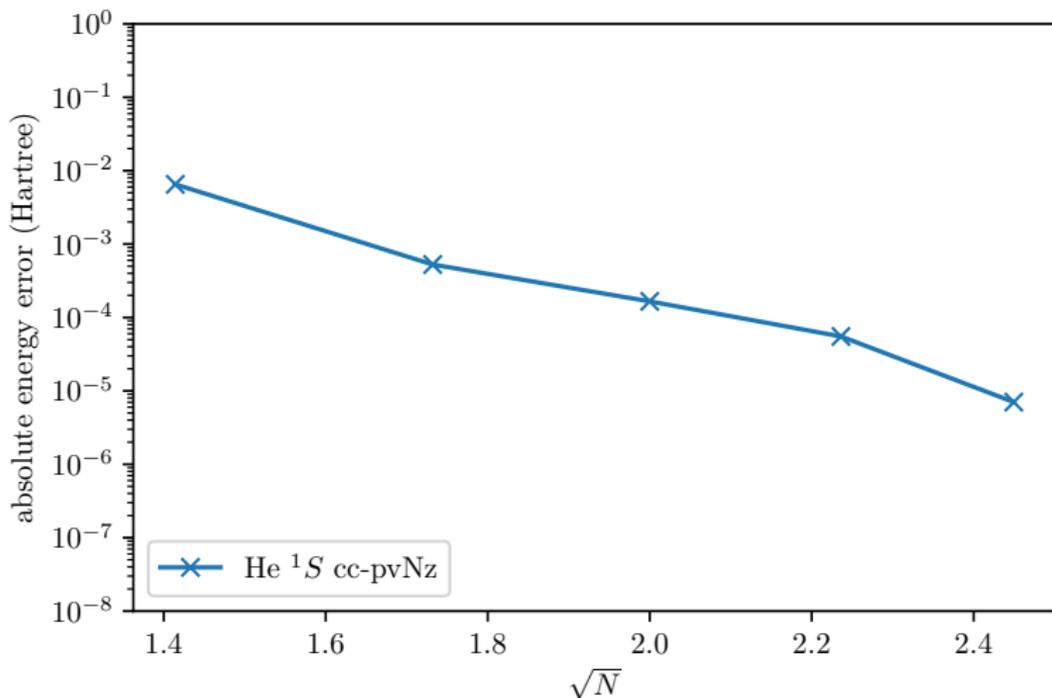
*Inria*



[https://michael-herbst.com/talks/2019.03.29\\_coulomb\\_sturmians\\_jussieu.pdf](https://michael-herbst.com/talks/2019.03.29_coulomb_sturmians_jussieu.pdf)

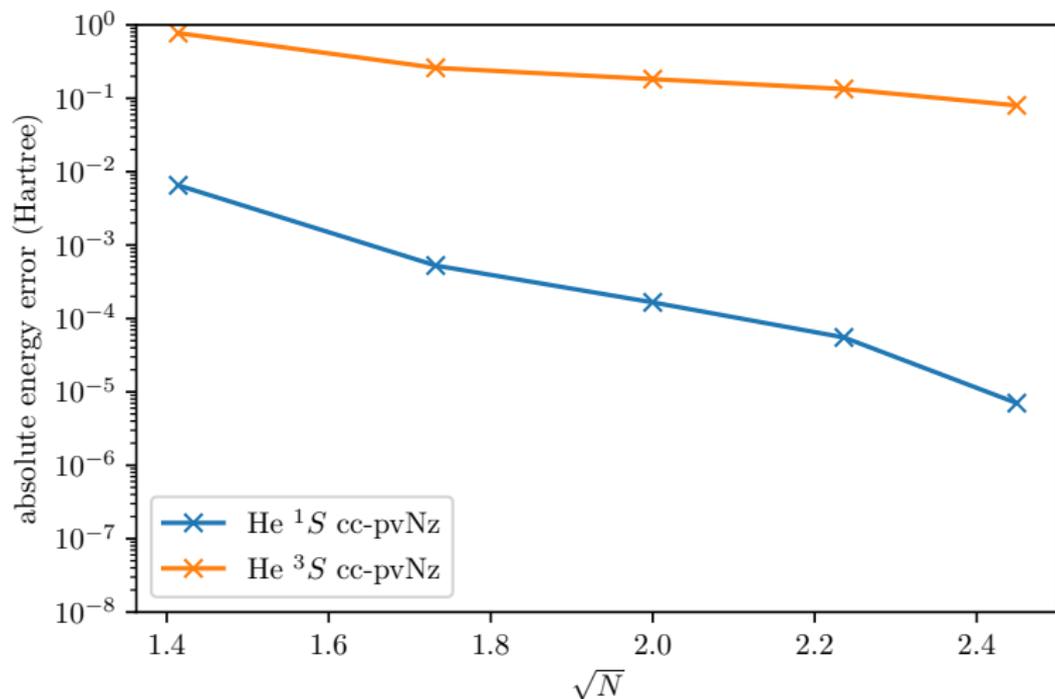
# The electronic structure of Helium

- Should be easy ... right?



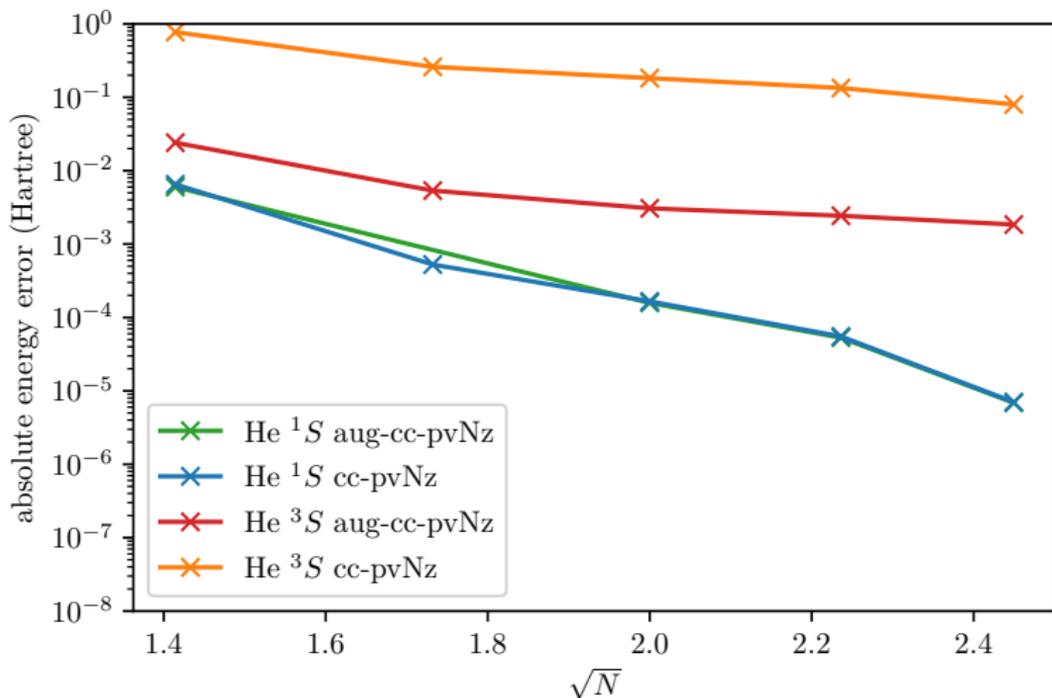
# The electronic structure of Helium

## ● Ups ...



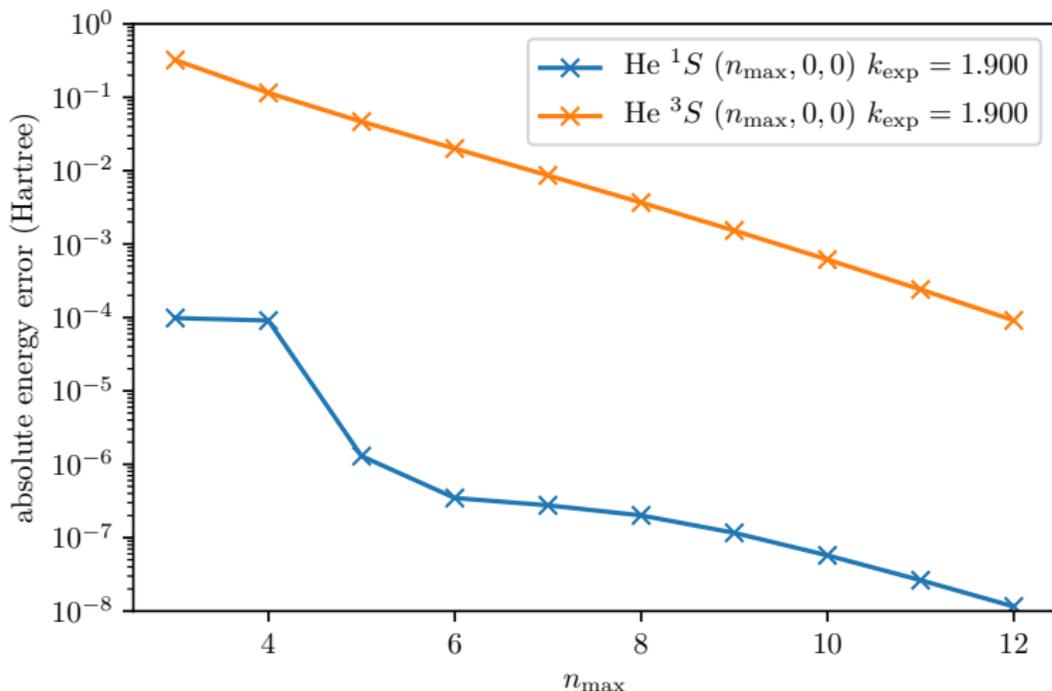
# The electronic structure of Helium

## ● Ahh ... Augmentation!

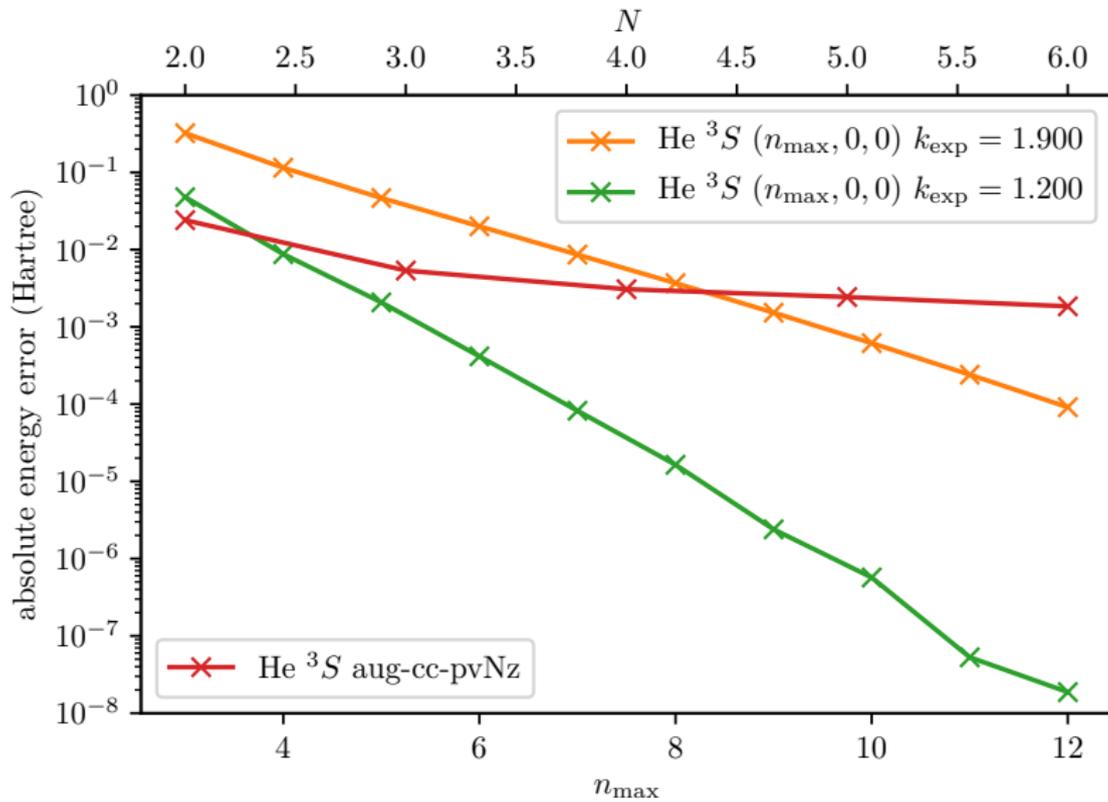


# The electronic structure of Helium

- Let's try ... Coulomb Sturmians



## Helium triplet ... direct comparison



## Beyond Gaussian-type orbitals

- Issues of cGTOs:
  - Representing the core region?
  - Representing the exponential decay?
  - Error estimates?
  - Black box modelling?
  - Distributed memory parallelisation?

⇒ Playing field to try other basis function types

# The potential of Coulomb Sturmians

- Prospect of CS:
  - Exponential-type basis functions
  - ✓ Represent the core region
  - ✓ Represent exponential decay
  - ? Error estimates
  - ? Black box modelling
  - ?? Distributed memory parallelisation

⇒ Implementation inside `molsturm`<sup>1</sup> and `sturmint`

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<sup>1</sup>M. F. Herbst, A. Dreuw and J. E. Avery. J. Chem. Phys., **149**, 84106 (2018)

# Contents

- 1 Properties of Coulomb Sturmian functions
- 2 Coulomb-Sturmian HF convergence studies
- 3 Ongoing work and speculative outlook
  - Helium  $^3S$  state
  - Correlated quantum-chemical methods
  - Modelling continuum-like states

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## Coulomb Sturmians

- Iso-energetic solutions  $\varphi_{nlm}$  to hydrogen-like equation<sup>1</sup>

$$\left(-\frac{1}{2}\Delta - \beta_n \frac{Z}{r}\right) \varphi_{nlm}(\mathbf{r}) = E \varphi_{nlm}(\mathbf{r})$$

- Scaling factor  $\beta_n$  chosen to uniform energy:

$$\beta_n = \frac{kn}{Z} \quad \Rightarrow \quad E = -\frac{k^2}{2}$$

- $\varphi_{nlm}$  look like hydrogenic orbitals with  $\frac{Z}{n}$  replaced by  $k$
- Radial part  $R_{nl}$  satisfies

$$\left(-\frac{1}{2r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r}\right) + \frac{l(l+1)}{2r^2} - \frac{nk}{r} - E\right) R_{nl}(r) = 0.$$

⇒ Sturm-Liouville equation<sup>2</sup>

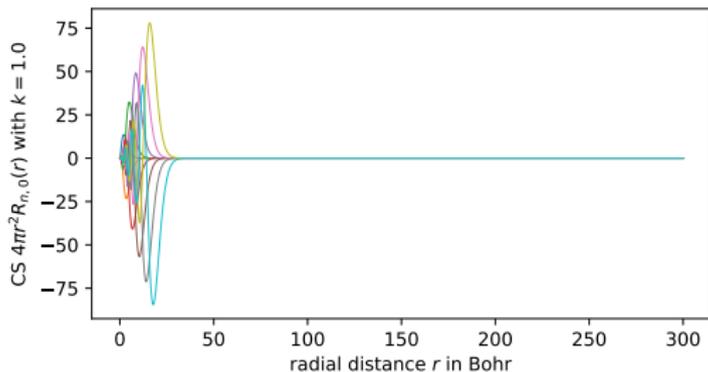
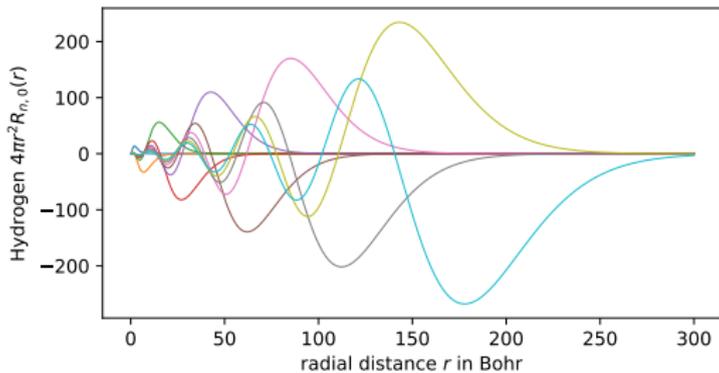
<sup>1</sup>H. Shull and P.-O. Löwdin. J. Chem. Phys., **30**, 617 (1959)

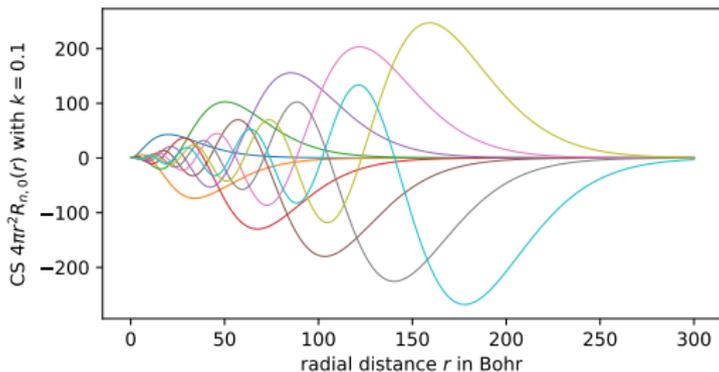
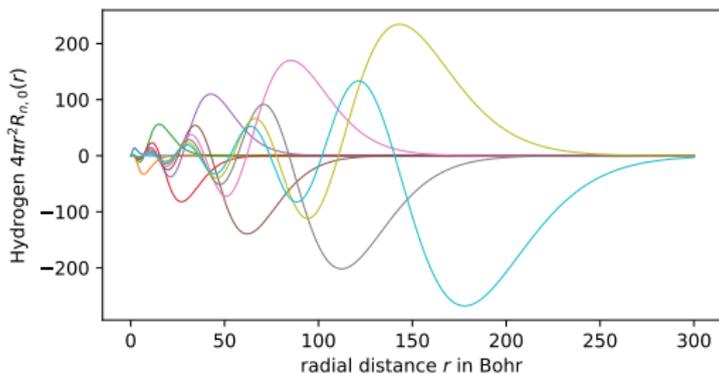
<sup>2</sup>M. Rotenberg. Ann. Phys., **19**, 262 (1962)

## Examples of the Coulomb Sturmian radial parts

$$\varphi_{\mu}(\underline{\mathbf{r}}) \equiv \varphi_{nlm}(\underline{\mathbf{r}}) = R_{nl}(r)Y_l^m(\hat{\mathbf{r}})$$

	$n$	$l$	$R_{nl}(r)$		
1s	1	0	$2k^{3/2}$		$\exp(-kr)$
2s	2	0	$2k^{3/2}$	$(1 - kr)$	$\exp(-kr)$
2p	2	1	$2k^{3/2}$	$kr$	$\exp(-kr)$
3s	3	0	$2k^{3/2}$	$\left(1 - 2kr + \frac{2}{3}(kr)^2\right)$	$\exp(-kr)$
3p	3	1	$2k^{3/2}$	$\frac{2\sqrt{2}}{3}kr \left(1 - \frac{kr}{2}\right)$	$\exp(-kr)$
3d	3	2	$2k^{3/2}$	$\frac{\sqrt{2}}{3\sqrt{5}}(kr)^2$	$\exp(-kr)$
general			$2k^{3/2}$	$P_{nl}(kr)$	$\exp(-kr)$

Hydrogen vs. CS radial distribution functions ( $k = 1.0$ )

Hydrogen vs. CS radial distribution functions ( $k = 0.1$ )

## Classical turning point and concentration

- Classical turning point: Zero kinetic energy
- Hydrogen-like functions:

$$-\frac{1}{n^2} = -\frac{Z}{r_T} \Leftrightarrow r_T = Zn^2$$

- Coulomb Sturmians:

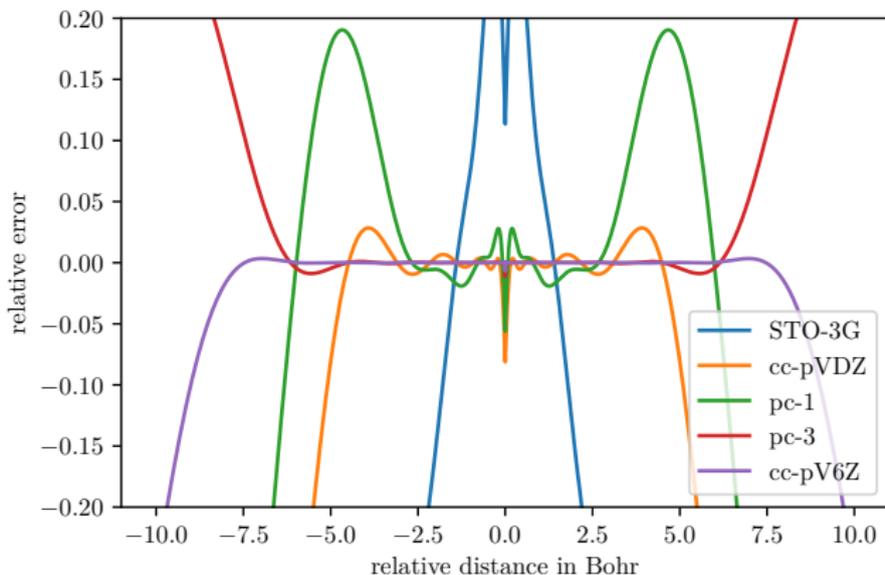
$$-\frac{k^2}{2} = -\beta_n \frac{Z}{r_T} \Leftrightarrow r_T = \frac{2Z\beta_n}{k^2} = \frac{2n}{k}$$

⇒ Hydrogen-like functions escape too fast into continuum

⇒ Exponent  $k$  allows to tune concentration of CS

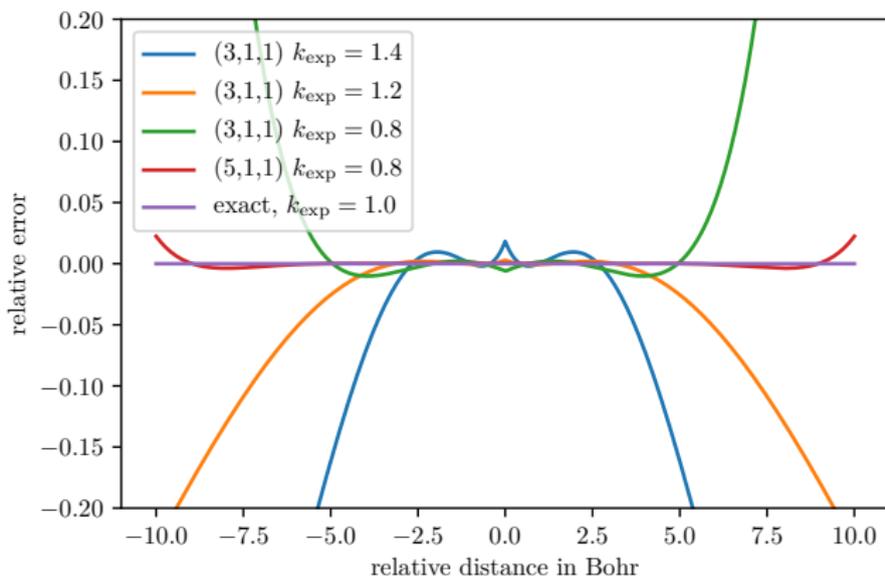
# Coulomb-Sturmians vs Gaussians (for hydrogen)

$$\varphi_{\mu}^{\text{GTO}}(\underline{\mathbf{r}}) = r^{l_{\mu}} \sum_i^{N_{\text{contr}}} c_{\mu,i} \exp(-\alpha_{\mu,i} r^2) \cdot Y_{l_{\mu}}^{m_{\mu}}(\hat{\mathbf{r}})$$



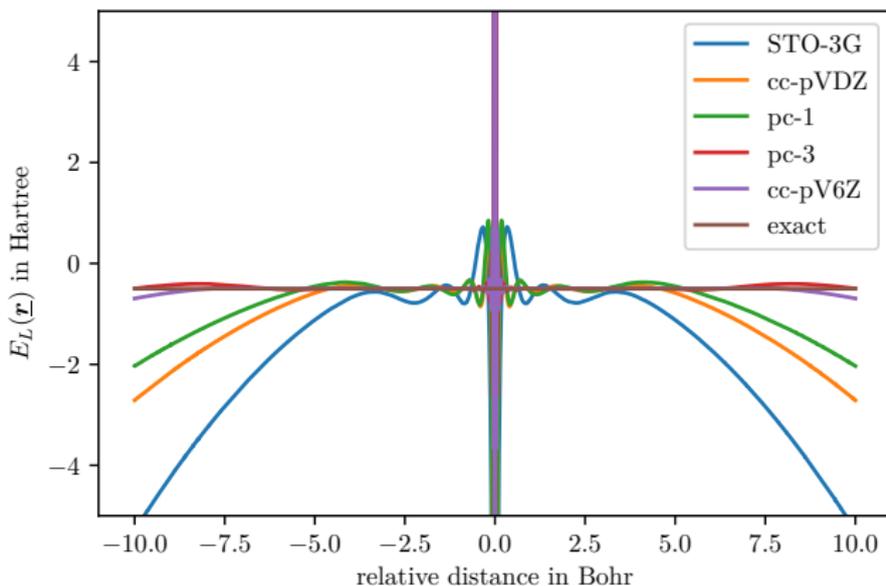
## Coulomb-Sturmians vs Gaussians (for hydrogen)

$$\varphi_{\mu}^{\text{CS}}(\underline{\mathbf{r}}) = 2k^{3/2} P_{nl}(2kr) \exp(-kr) \cdot Y_l^m(\hat{\mathbf{r}})$$



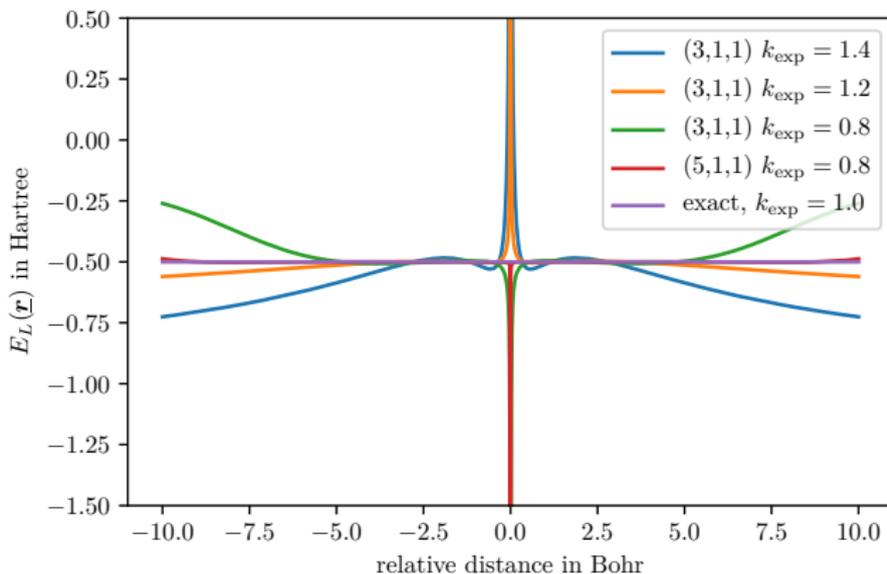
## Coulomb-Sturmians vs Gaussians (for hydrogen)

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## Coulomb-Sturmians vs Gaussians (for hydrogen)

$$\varphi_{\mu}^{\text{CS}}(\underline{\mathbf{r}}) = 2k^{3/2} P_{nl}(2kr) \exp(-kr) \cdot Y_l^m(\hat{\mathbf{r}})$$



## (Atomic) Coulomb-Sturmians properties overview

- Complete (Schauder) basis<sup>12</sup>
- Correct physics: Nuclear cusp and exponential decay
- CS exponent  $k$ :
  - Identical for each basis function
  - Controls concentration of basis functions
- One-electron integrals sparse and analytic
- Two-electron integrals sparse tensor contraction
- Generalisation to molecules possible<sup>3</sup>

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<sup>1</sup>B. Klahn and W. A. Bingel. *Theoretica chimica acta*, **44**, 9 (1977)

<sup>2</sup>B. Klahn and W. A. Bingel. *Theoretica chimica acta*, **44**, 27 (1977)

<sup>3</sup>J. E. Avery and J. S. Avery. vol. 70 of *Advances in Quantum Chemistry*, 265 – 324. Academic Press (2015)



## Evaluating integrals: Useful relations

- Potential-weighted orthonormality relation:

$$\int_{\mathbb{R}^3} \varphi_{\mu'}^*(\underline{\mathbf{r}}) \frac{n}{kr} \varphi_{\mu}(\underline{\mathbf{r}}) d\underline{\mathbf{r}} = \delta_{\mu\mu'}$$

- Fourier-transform of Coulomb Sturmians:

$$\begin{aligned} \hat{\varphi}_{\mu}(\underline{\mathbf{p}}) &= \frac{1}{(2\pi)^{3/2}} \int_{\mathbb{R}^3} \varphi_{\mu}(\underline{\mathbf{r}}) e^{-i\underline{\mathbf{p}} \cdot \underline{\mathbf{x}}} d\underline{\mathbf{r}} \\ &= \frac{4k^{5/2}}{(k^2 + p^2)^2} Y_{\lambda,l,m}(\hat{\underline{\mathbf{u}}}) \end{aligned}$$

with  $\underline{\mathbf{p}} \equiv (p_1, p_2, p_3)^T$ ,  $\lambda \equiv n - 1$  and

$$\hat{\underline{\mathbf{u}}} \equiv \left( \frac{2kp_1}{k^2 + p^2}, \frac{1kp_2}{k^2 + p^2}, \frac{2kp_3}{k^2 + p^2}, \frac{k^2 - p^2}{k^2 + p^2} \right)^T$$

- **Fock transformation** to hyperspherical harmonics  $Y_{\lambda,l,m}(\underline{\mathbf{u}})$ <sup>1</sup>

<sup>1</sup>V. Fock. Zeitschrift für Physik, **98**, 145 (1935)

# One-electron integrals

- Overlap:

$$S_{\mu',\mu} = \delta_{m',m} \delta_{l',l} \begin{cases} n = n' & 1 \\ |n - n'| = 1 & s_{n,n+1}^{(l)} \in \mathbb{R} \\ \text{else} & 0 \end{cases}$$

where

$$s_{n,n+1}^{(l)} = -\frac{1}{2} \sqrt{\frac{(n-l)(n+l+1)}{n(n+1)}}$$

- Nuclear attraction:

$$V_{\mu',\mu} = -\delta_{\mu',\mu} \frac{kZ}{n}$$

- Kinetic:

$$T_{\mu',\mu} = k^2 (\delta_{\mu',\mu} - S_{\mu',\mu})$$

## Two-electron integrals

- ERI tensor from kernel  $I_{\mu\nu}$

$$(\varphi_{\mu_1}\varphi_{\mu_2}|\varphi_{\mu_3}\varphi_{\mu_4}) = \sum_{\mu\nu} C_{\mu_2\mu_1}^{\mu} I_{\mu\nu} C_{\mu_3\mu_4}^{\nu}$$

where

$$I_{\mu\mu'} = \delta_{ll'}\delta_{mm'}I_{nn'}^{(l)},$$

thus sparse and storable (30 kB for  $n = 20$ ).

- The  $C_{\mu_1\mu_2}^{\mu}$  are the coefficients in

$$\varphi_{\mu_1}^*(\mathbf{r})\varphi_{\mu_2}(\mathbf{r}) = \sum_{\mu} C_{\mu_1\mu_2}^{\mu}\varphi_{\mu}(2k, \mathbf{r}),$$

again sparse (220 MB for  $n = 20$ ).

⇒ Like *exact* density fitting

Evaluating molecular CS 1e integrals<sup>1</sup>

- Let  $\tau \equiv (A, n, l, m)$  and  $\varphi_\tau(\underline{\mathbf{r}}) \equiv \varphi_\mu(\underline{\mathbf{r}} - \underline{\mathbf{R}}_A)$
- Key quantity: **Shibuya-Wulfman integrals**

$$\begin{aligned} \mathfrak{S}_{\tau\tau'} &= \int_{\mathbb{R}^3} \varphi_\tau^*(\underline{\mathbf{r}}) \left( \frac{-\Delta + k^2}{2k^2} \right) \varphi_{\tau'}(\underline{\mathbf{r}}) d\underline{\mathbf{r}} \\ &= (2\pi)^{3/2} \sum_{\mu''} f_{n'',l''}(S) Y_{l''}^{m''}(\hat{\underline{\mathbf{S}}}) c_{\mu'';\mu,\mu'} \end{aligned}$$

with  $\underline{\mathbf{S}} = k(\underline{\mathbf{R}}_A - \underline{\mathbf{R}}_{A'})$ ,  $c_{\mu'';\mu,\mu'}$  sparse and pre-computable

- $f_{nl}(S)$ : analytic in  $S$ , related to CS radial part  $R_{nl}(S)$

<sup>1</sup>J. E. Avery and J. S. Avery. vol. 70 of *Advances in Quantum Chemistry*, 265 – 324. Academic Press (2015)

# Evaluating molecular CS ERI integrals<sup>1</sup>

- Following a similar idea as before

$$(\varphi_{\tau_1} \varphi_{\tau_2} | \varphi_{\tau_3} \varphi_{\tau_4}) \simeq \sum_{\mu\nu} C_{\tau_2\tau_1}^\tau I_{\tau\tau'} C_{\tau_3\tau_4}^{\tau'}$$

where

$$\varphi_{\tau_1}(\underline{\mathbf{r}}) \varphi_{\tau_2}^*(\underline{\mathbf{r}}) \simeq \sum_{\tau} C_{\tau_2\tau_1}^\tau \varphi_{\tau}(2k, \underline{\mathbf{r}})$$

with  $\sum_{\tau} \equiv \sum_{A_1\mu} + \sum_{A_2\mu}$ .

- For this: Solve **once** linear system (thus  $\simeq$ ):

$$\sum_{\tau'} S_{\tau\tau'}(2k, \underline{\mathbf{S}}) C_{\tau_2\tau_1}^{\tau'} = \int_{\mathbb{R}^3} \varphi_{\tau}^*(2k, \underline{\mathbf{r}}) \varphi_{\tau_1}(\underline{\mathbf{r}}) \varphi_{\tau_2}^*(\underline{\mathbf{r}}) d\underline{\mathbf{r}}$$

<sup>1</sup>J. E. Avery and J. S. Avery. vol. 70 of *Advances in Quantum Chemistry*, 265 – 324. Academic Press (2015)

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## Construction of CS basis sets

- All CS basis functions share common exponent  $k$
- Index restrictions by CS equations:

$$n > 0 \qquad 0 \leq l < n \qquad -l \leq m \leq l$$

- Basis set construction:<sup>1</sup> Additional restrictions

$$n \leq n_{\max} \qquad l \leq l_{\max} \qquad m \leq m_{\max}$$

- Restriction by  $m_{\max}$ , sometimes  $l_{\max}$  neglected
- ⇒ Physically motivated

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<sup>1</sup>M. F. Herbst, J. E. Avery and A. Dreuw. Physical Review A, **99**, 012512 (2019)

## The parameters $n_{\max}$ , $l_{\max}$ and $m_{\max}$

- Sturm-Liouville form of radial part (for each  $l$ ):

$$\forall 0 \leq l < n : R_{nl} \in \text{span}\{R_{n',0}\}_{n' \leq n}$$

⇒  $n_{\max}$  controls radial discretisation

⇒  $l_{\max}$  and  $m_{\max}$  control angular discretisation

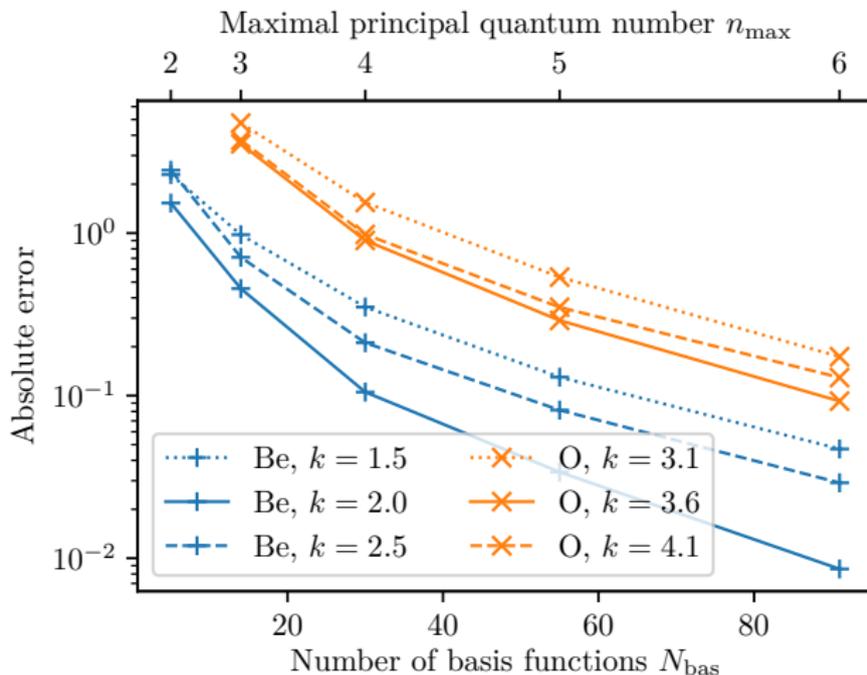
- Scaling of basis size if only restricted by  $n_{\max}$ :

$$N_{\text{bas}}(n_{\max}) = \frac{(2n_{\max} + 1)(n_{\max} + 1)n_{\max}}{6} \in \mathcal{O}(n_{\max}^3)$$

- Scaling of basis size if restricted by  $n_{\max}$  and  $l_{\max}$ :

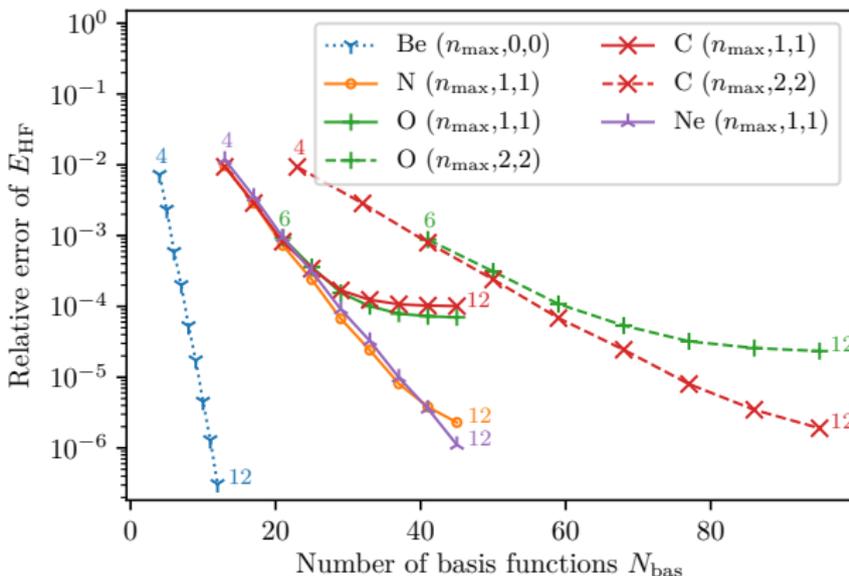
$$N_{\text{bas}}(n_{\max}) = (n_{\max} - 1)(l_{\max} + 1)^2 \in \mathcal{O}(n_{\max} l_{\max}^2)$$

# Energy convergence<sup>1</sup>

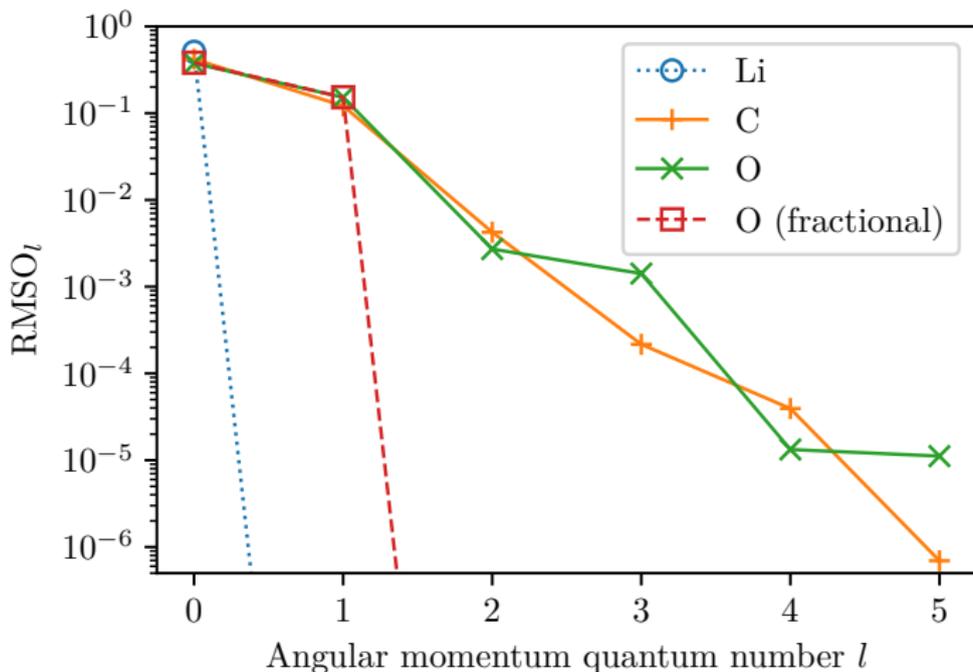


<sup>1</sup>M. F. Herbst, J. E. Avery and A. Dreuw. Physical Review A, **99**, 012512 (2019)

# Basis progressions: Energy versus basis size<sup>1</sup>



<sup>1</sup>M. F. Herbst, J. E. Avery and A. Dreuw. Physical Review A, **99**, 012512 (2019)

Energy convergence: Deviating behaviour of elements<sup>1</sup>

<sup>1</sup>M. F. Herbst, J. E. Avery and A. Dreuw. Physical Review A, **99**, 012512 (2019)

## Root mean square coefficient value per $l$

- Root mean square occupied coefficient par AM  $l$ :

$$\text{RMSO}_l = \sqrt{\sum_{\mu} \sum_i \sum_{\sigma \in \{\alpha, \beta\}} \frac{1}{N_{\text{elec}}^{\sigma} N_{\text{bas}, l}} \left( C_{\mu, i}^{\sigma} f_{ii}^{\sigma} \right)^2}$$

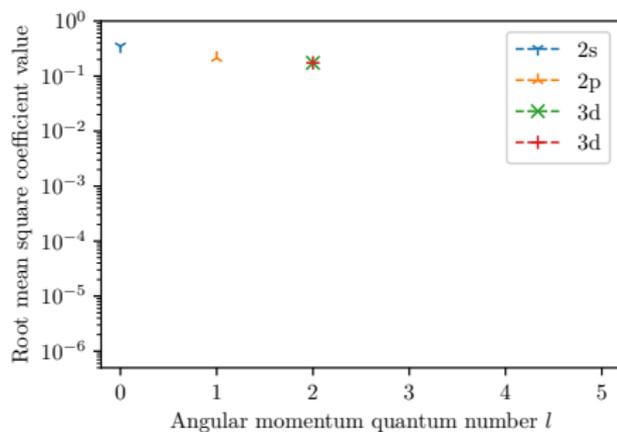
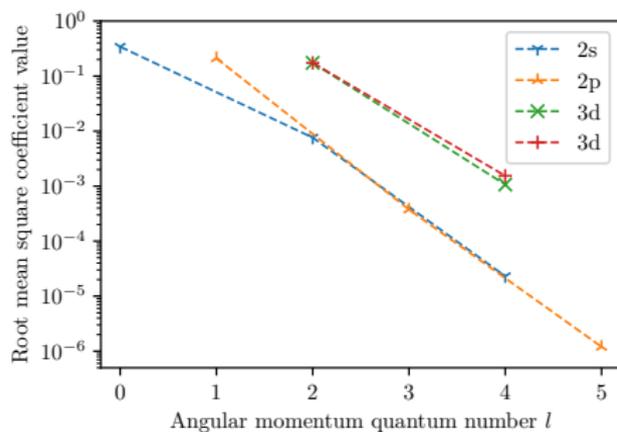
where  $i$  runs over orbitals,  $\mu$  over basis functions and

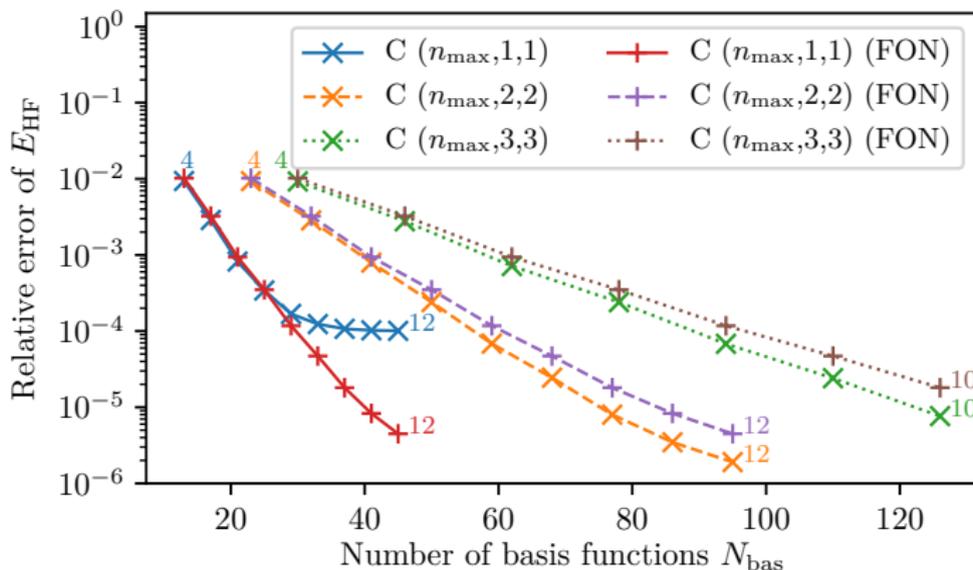
$N_{\text{elec}}^{\sigma}$  Number of electrons with spin  $\sigma$

$f_{ii}^{\sigma}$  Occupation number

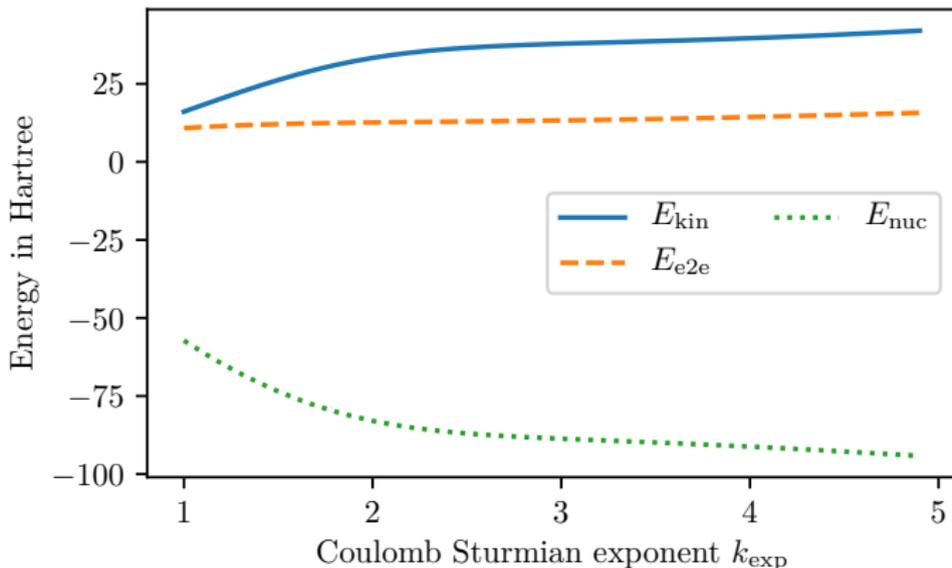
$N_{\text{bas}, l}$  Number of basis functions of AM  $l$

- Measure for importance of  $l$  in HF wave function

RMSO<sub>l</sub> plots<sup>1</sup><sup>1</sup>M. F. Herbst, J. E. Avery and A. Dreuw. Physical Review A, **99**, 012512 (2019)

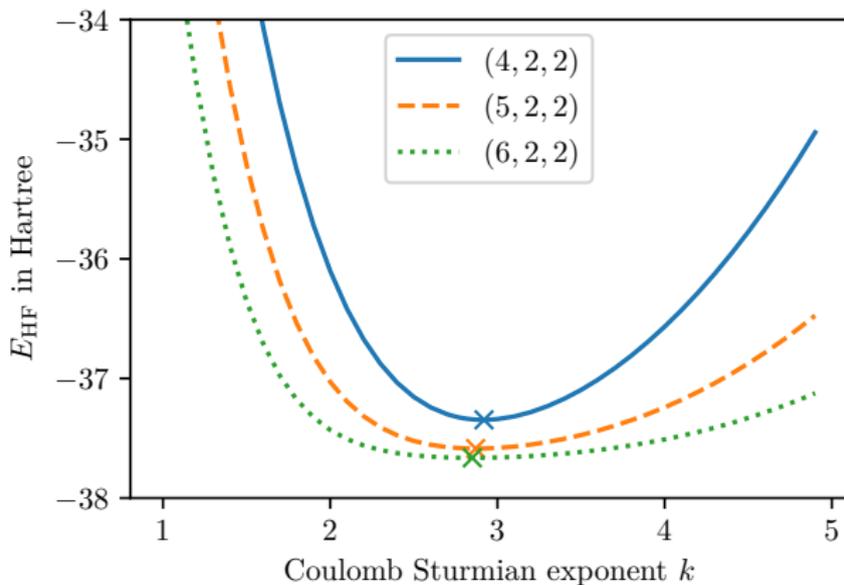
Basis progressions: Carbon<sup>1</sup><sup>1</sup>M. F. Herbst, J. E. Avery and A. Dreuw. Physical Review A, **99**, 012512 (2019)

# $k$ -dependence terms carbon<sup>1</sup>

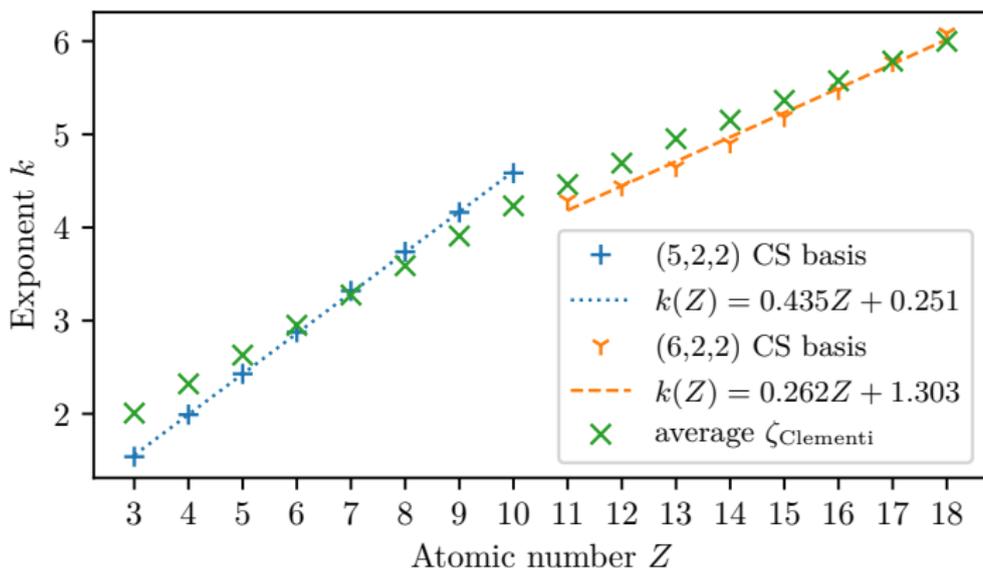


<sup>1</sup>M. F. Herbst, J. E. Avery and A. Dreuw. Physical Review A, **99**, 012512 (2019)

# $k$ -dependence terms carbon<sup>1</sup>



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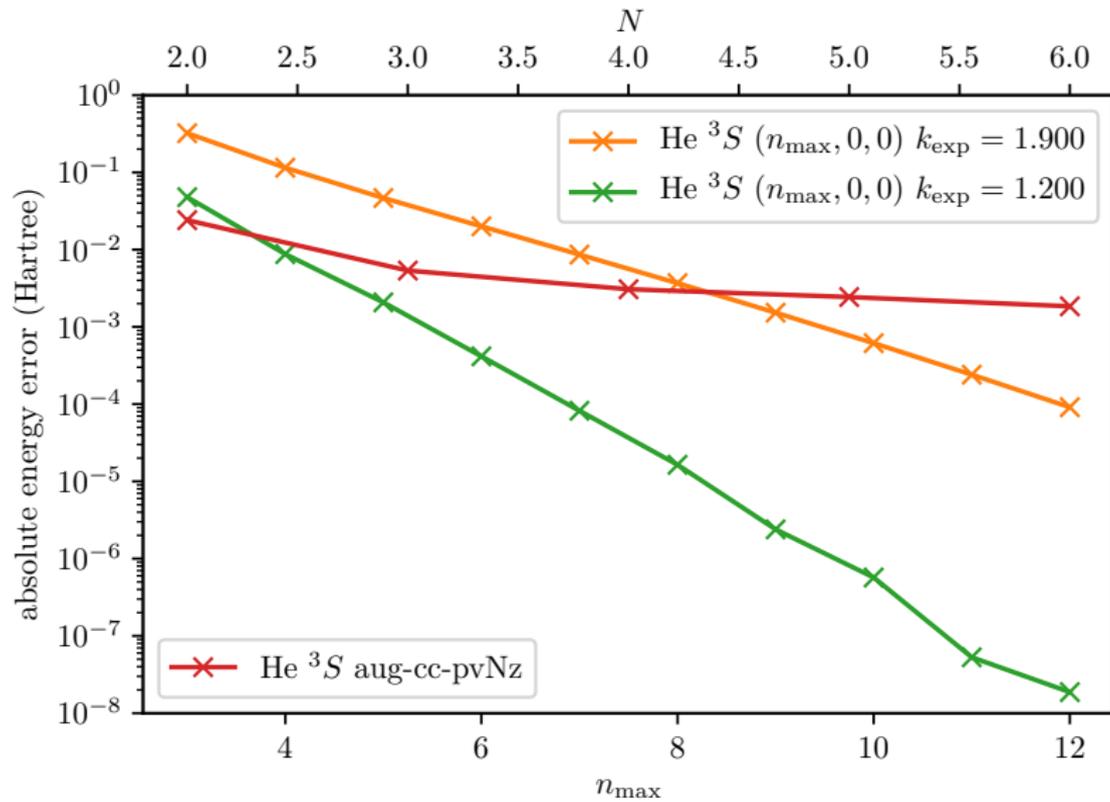
Optimal exponent versus Clementi exponents<sup>1</sup>

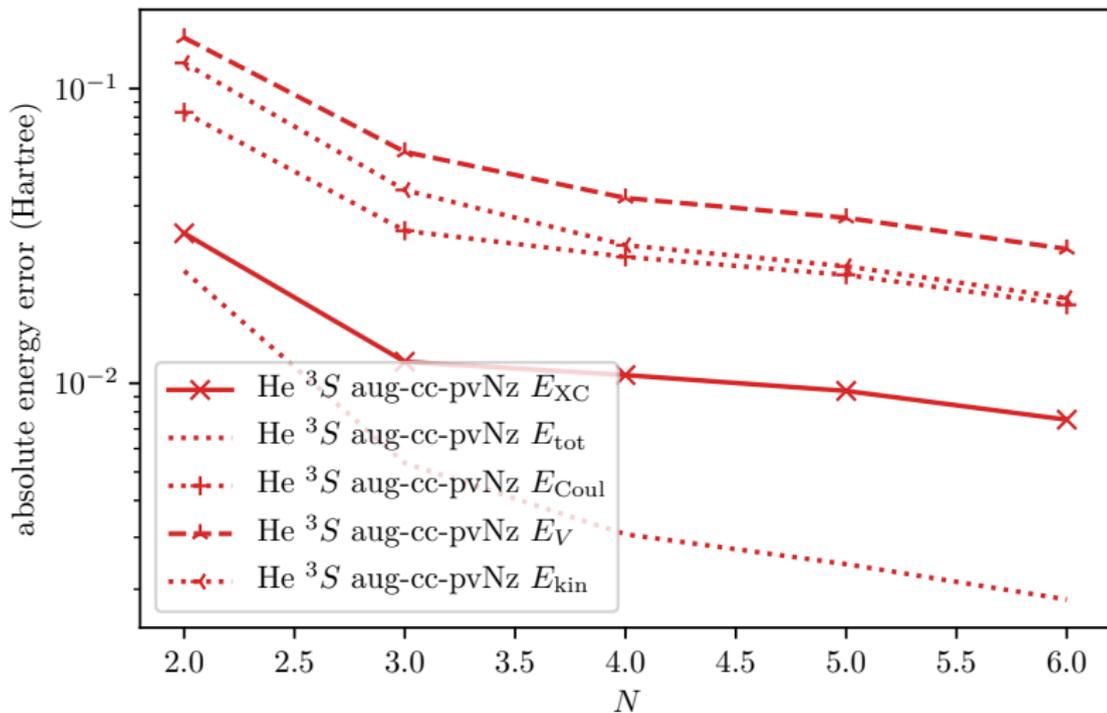
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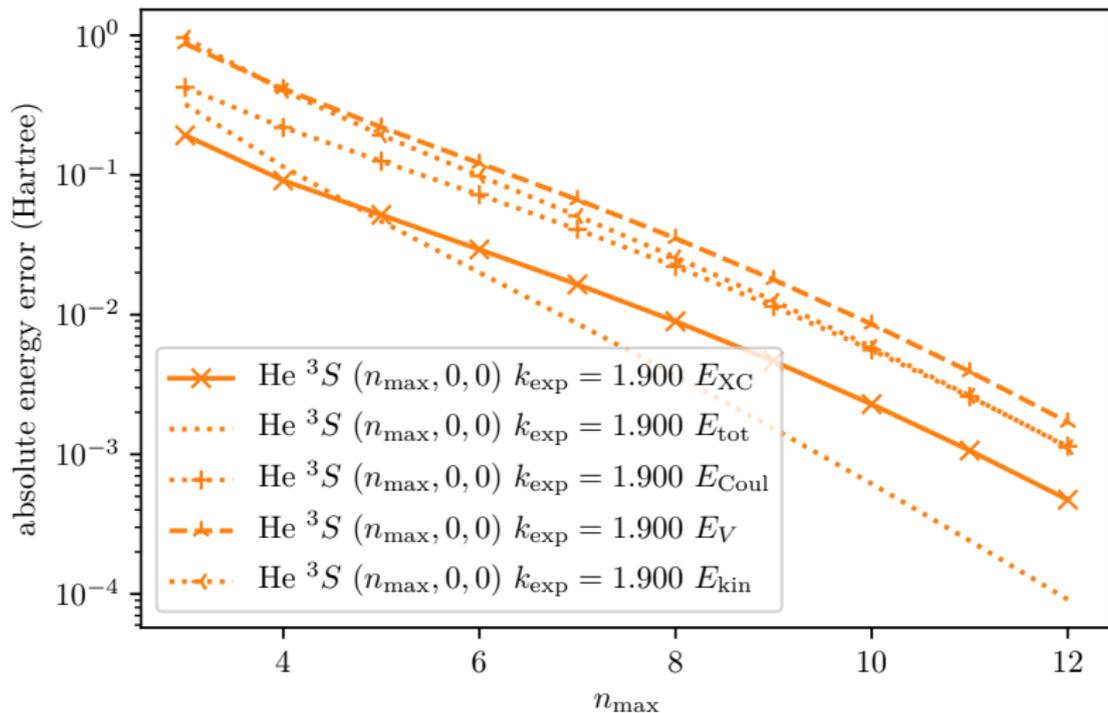
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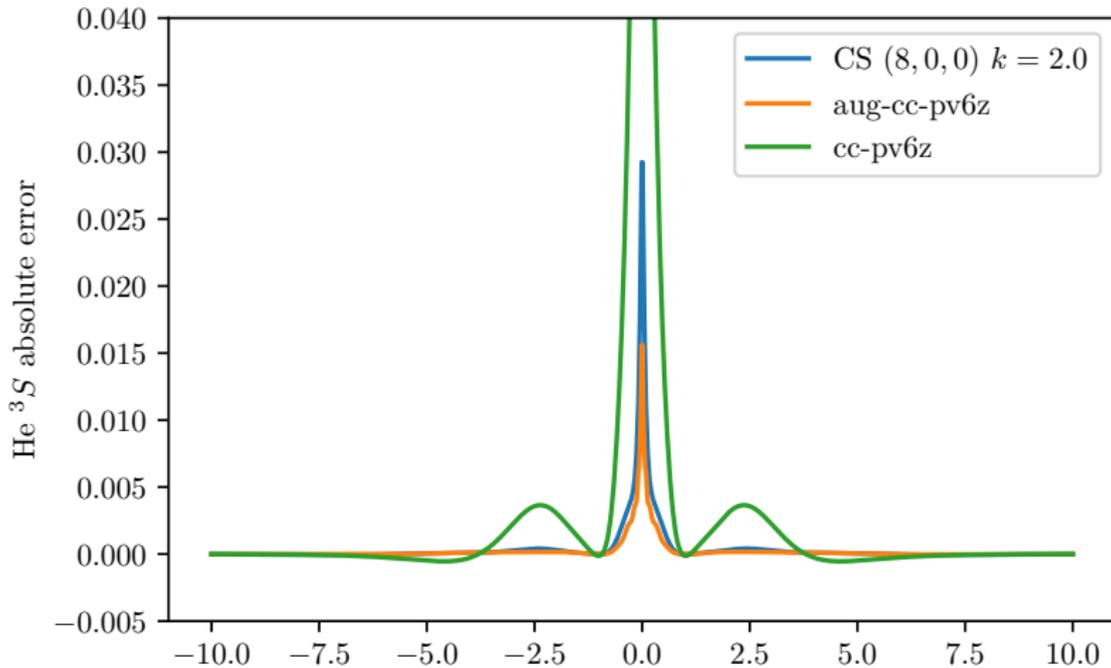
## Returning to our example from before ...



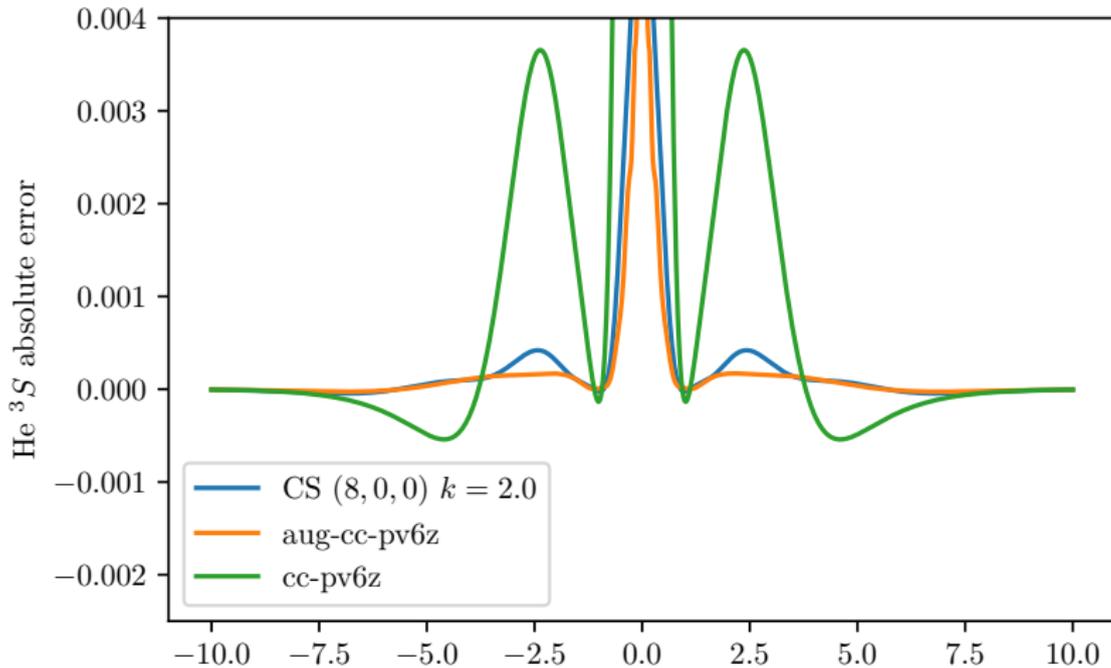
Helium  $^3S$  term-wise convergence

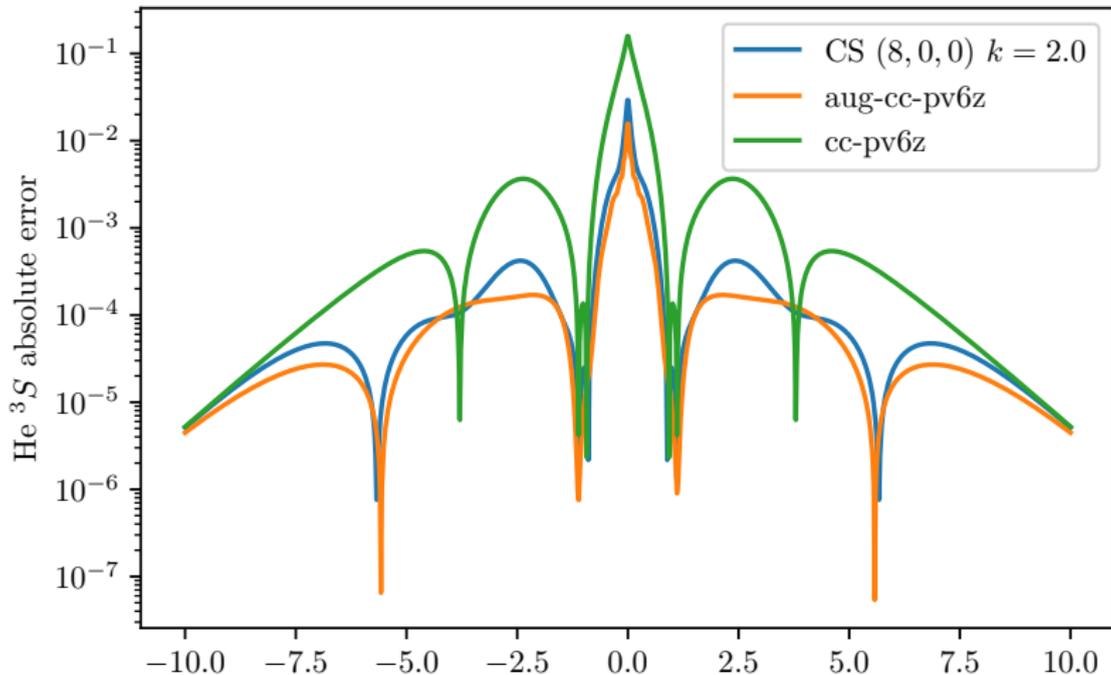
Helium  $^3S$  term-wise convergence

# Helium $^3S$ density error

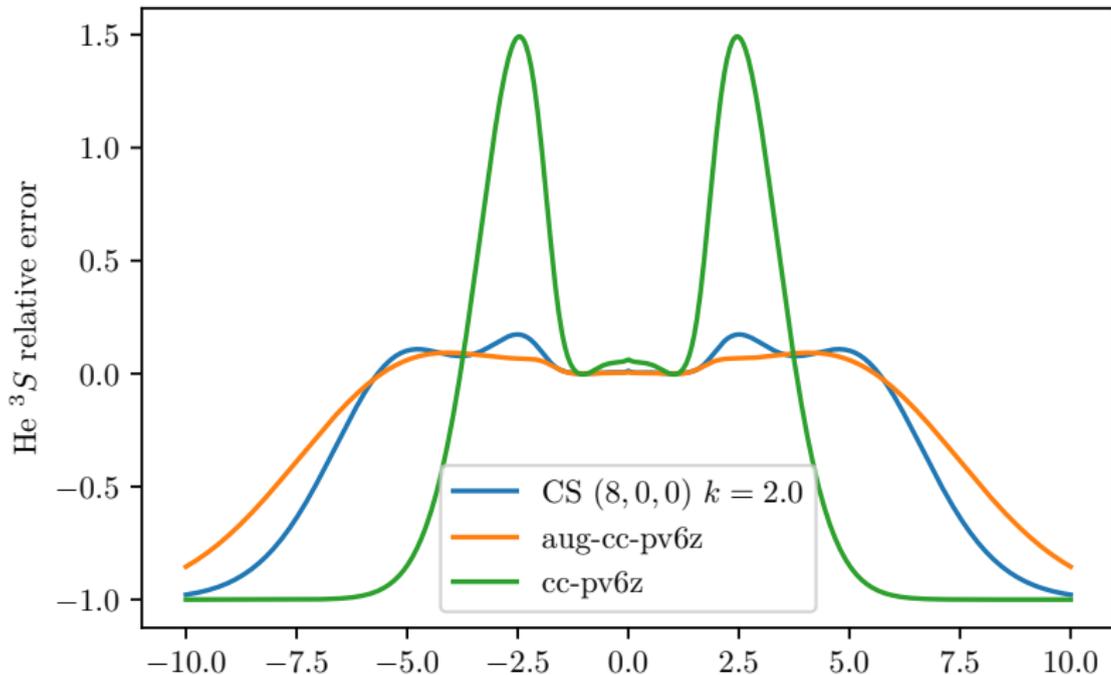


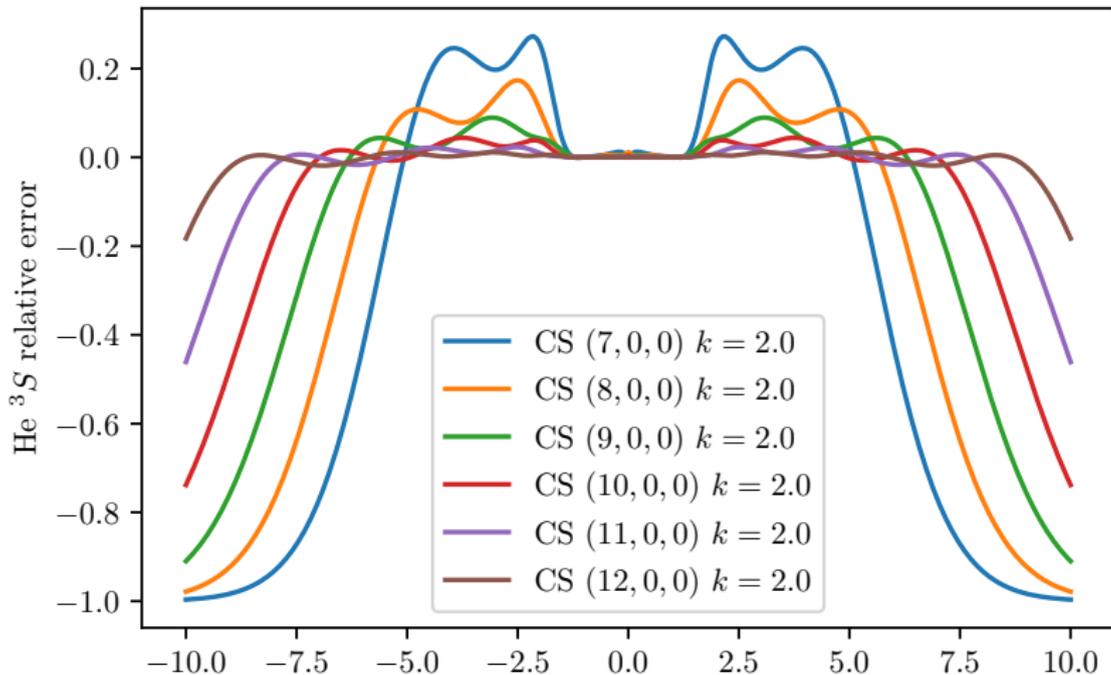
# Helium $^3S$ density error



Helium  $^3S$  density error

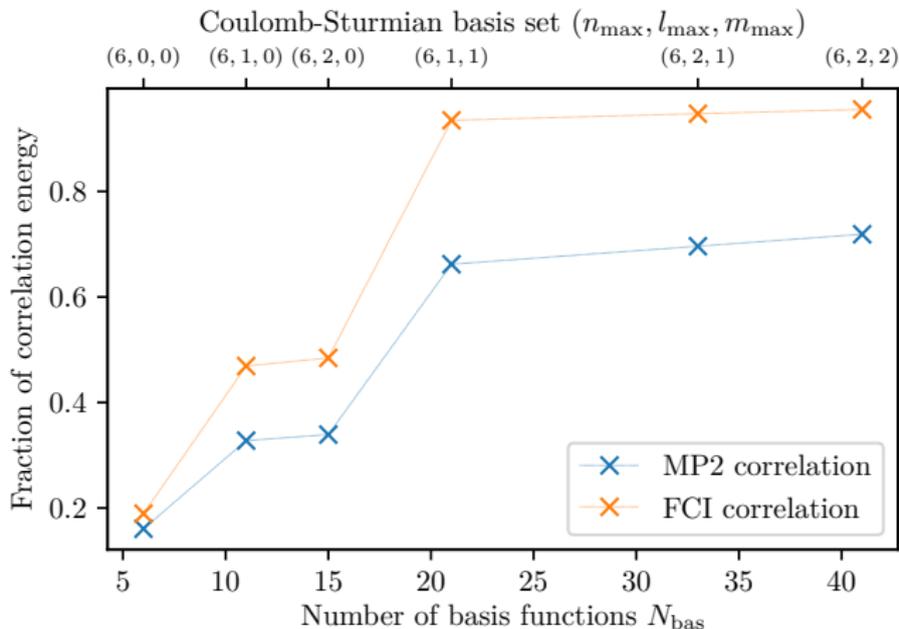
# Helium $^3S$ density error



Helium  $^3S$  density error

# MP(2) and FCI (basis size)<sup>1</sup>

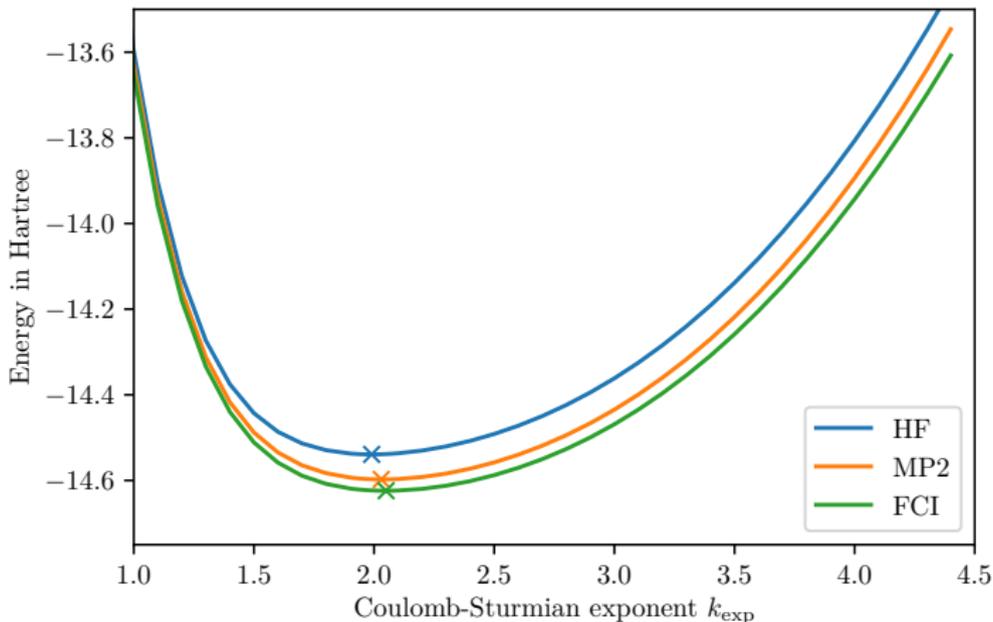
- Beryllium,  $k = 2.1$ , correlation energy relative to FCI (10, 2, 2)



<sup>1</sup>M. F. Herbst. Ph.D. thesis, Ruprecht-Karls-Universität Heidelberg (2018)

# MP(2) and FCI ( $k$ dependence)<sup>1</sup>

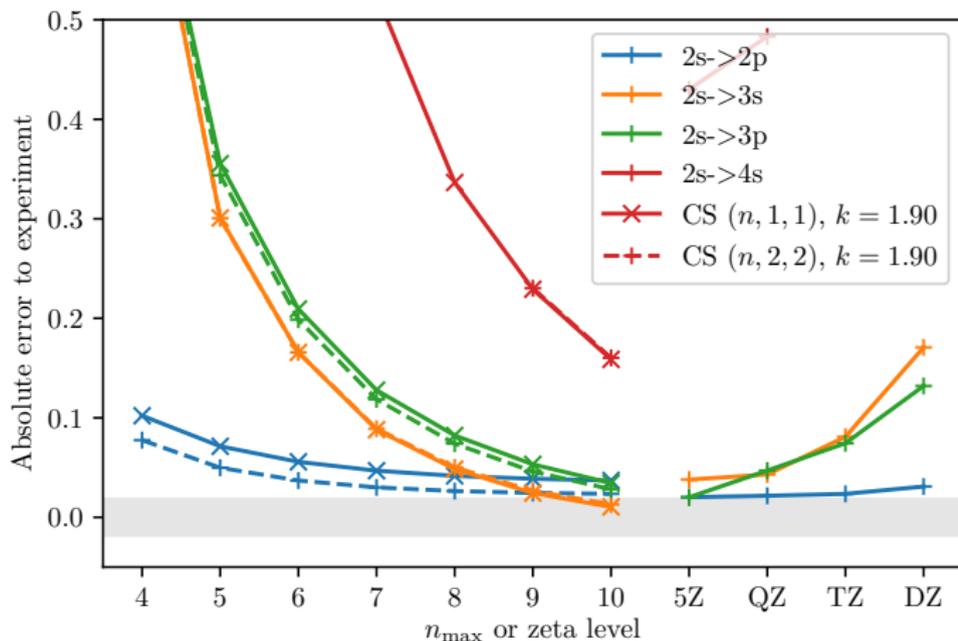
- Beryllium, (5, 1, 1) CS basis



<sup>1</sup>M. F. Herbst. Ph.D. thesis, Ruprecht-Karls-Universität Heidelberg (2018)

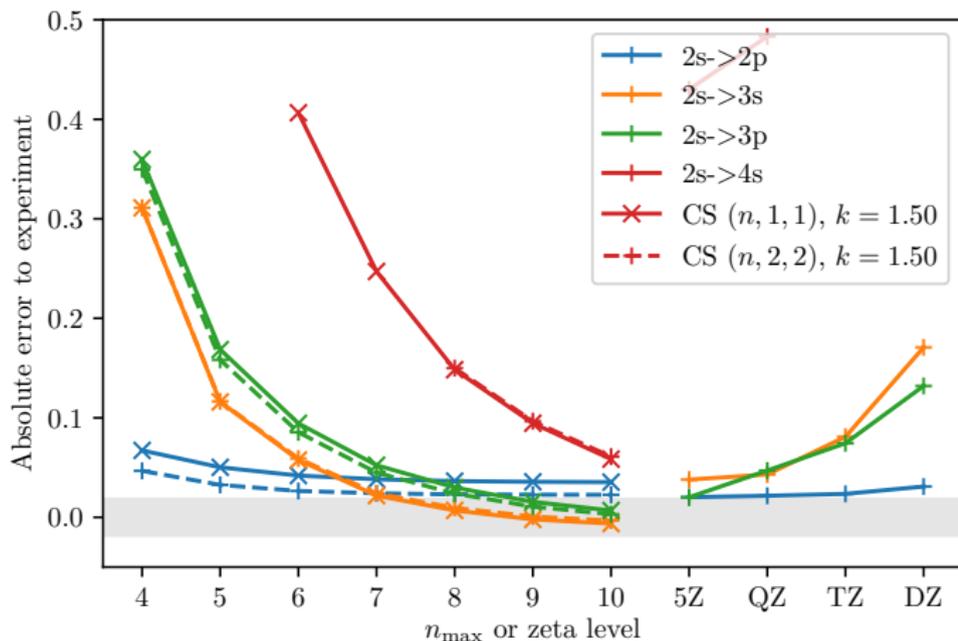
# ADC(2) absolute error

## ● Beryllium CS progressions vs. Gaussians



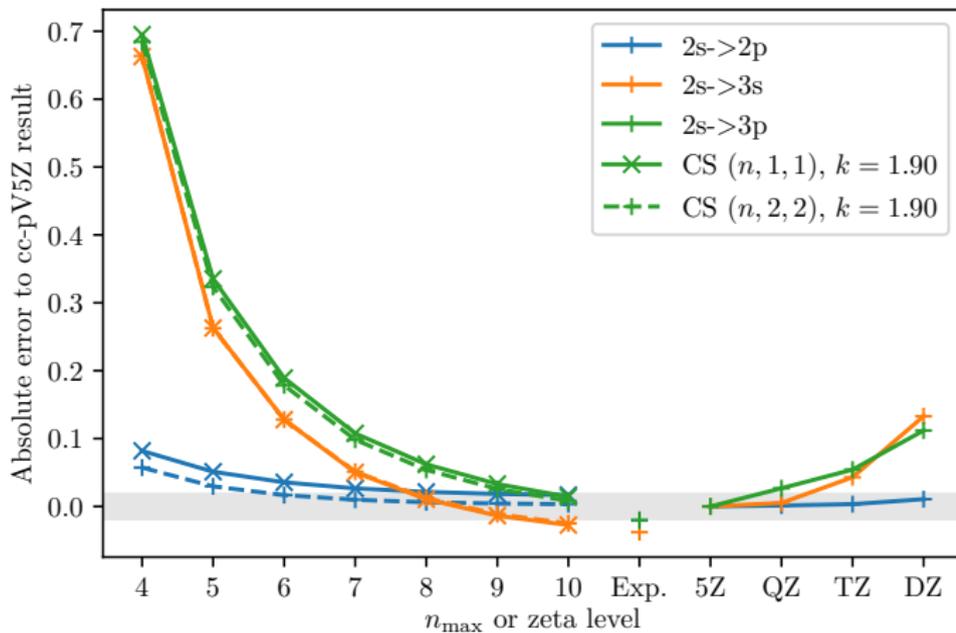
# ADC(2) absolute error

- Beryllium CS progressions vs. Gaussians



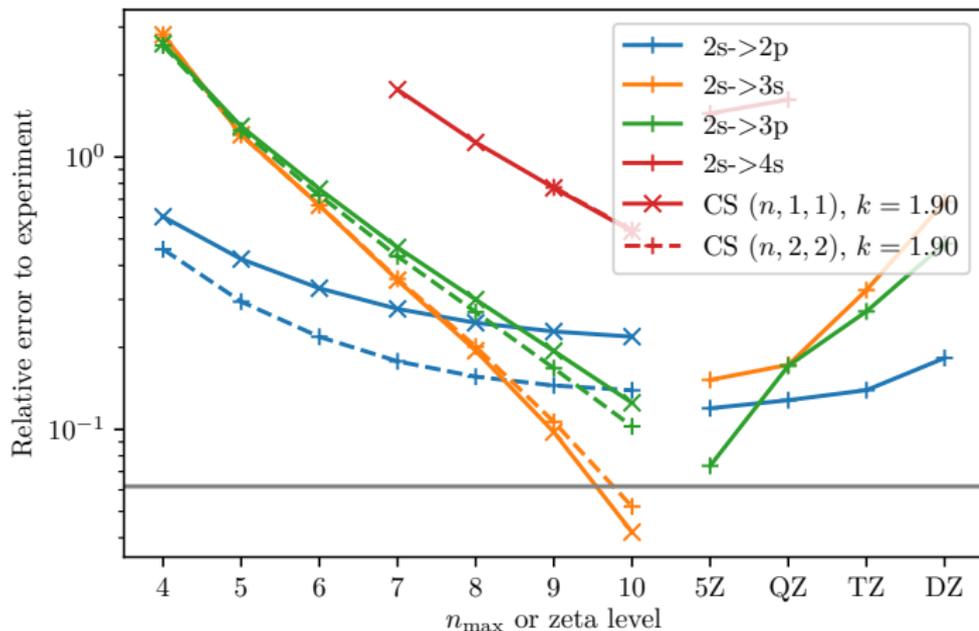
# ADC(2) Gaussian error

- Beryllium CS progressions vs. Gaussians



## ADC(2) relative error

## ● Beryllium CS progressions vs. Gaussians



## $k$ -targeted FCI for generalised Sturmians

- Generalised Sturmians: Atomic  $N$ -electron basis
  - Relation  $E = -k^2/2$  generalises to these Sturmians
  - $k$ -targeted formulation:
    - Instead of solving for  $E$ , solve directly for parameters  $k$
- ⇒ Each state has a different set of  $k$ s
- Similar formulation for Coulomb-Sturmian basis?
    - ⇒ Hints how to do this in an **atomic** CI-like context<sup>1</sup>
      - Not aware of any HF / CI work on this for a CS basis

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<sup>1</sup>J. E. Avery and J. S. Avery. *Generalized Sturmians and Atomic Spectra*.  
World Scientific (2006)

## Coulomb-Sturmians for correlated methods

- Convergence behaviour?
- Systematic basis set construction?
- Contracted Coulomb-Sturmian basis sets?
- Basis sets with multiple  $k$ s?
- Molecular Coulomb Sturmians?
- $k$ -targeted formulations for HF and configuration interaction?
- $k$ -targeted formulations for molecular problems?

## CS-based modelling of continuum-like states ?

- Atomic CS functions are a complete basis
- Parameters  $k$  can be tuned to capture features of interest
- $k$ -targeted formulation could automatise this
- Sturmian-type functions successfully employed to treat photoionisation processes <sup>1 2 3</sup>

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<sup>1</sup>D. M. Mitnik, F. D. Colavecchia, G. Gasaneo and J. M. Randazzo. *Comp. Phys. Comm.*, **182**, 1145 (2011)

<sup>2</sup>J. M. Randazzo, D. Mitnik, G. Gasaneo, L. U. Ancarani and F. D. Colavecchia. *Eur. Phys. J.D*, **69**, 189 (2015)

<sup>3</sup>C. M. Granados-Castro, L. U. Ancarani, G. Gasaneo and D. M. Mitnik. vol. 73 of *Advances in Quantum Chemistry*, 3 – 57. Academic Press (2016)

## Coulomb Sturmians in essence

- Complete, exponential-type basis
  - ⇒ May capture cusp and decay behaviour
  - ⇒ Apparently less care required to converge special systems
- Basis set construction for HF problems:<sup>1</sup>
  - Physical rationale for arising parameters
  - Convergence explained by physical arguments
  - Implementation for atoms: `sturmint` and `molsturm`<sup>2</sup>
- Extension to molecular integrals feasible
- Basis sets for correlated methods required

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<sup>1</sup>M. F. Herbst, J. E. Avery and A. Dreuw. Physical Review A, **99**, 012512 (2019)

<sup>2</sup>M. F. Herbst, A. Dreuw and J. E. Avery. J. Chem. Phys., **149**, 84106 (2018)

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# Questions?

**Papers:** M. F. Herbst, A. Dreuw and J. E. Avery. J. Chem. Phys., **149**, 84106 (2018)

M. F. Herbst, J. E. Avery and A. Dreuw. Physical Review A, **99**, 012512 (2019)

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# Achievements of molsturm

- Basis-function independent design
  - Plug and play new discretisations
  - Basis-type agnostic SCF procedure
- Easy-to-use interfaces
  - Integrate with existing code (e.g. Post-HF)
  - Avoid reinventing the wheel
  - Rapid prototyping, testing and analysis

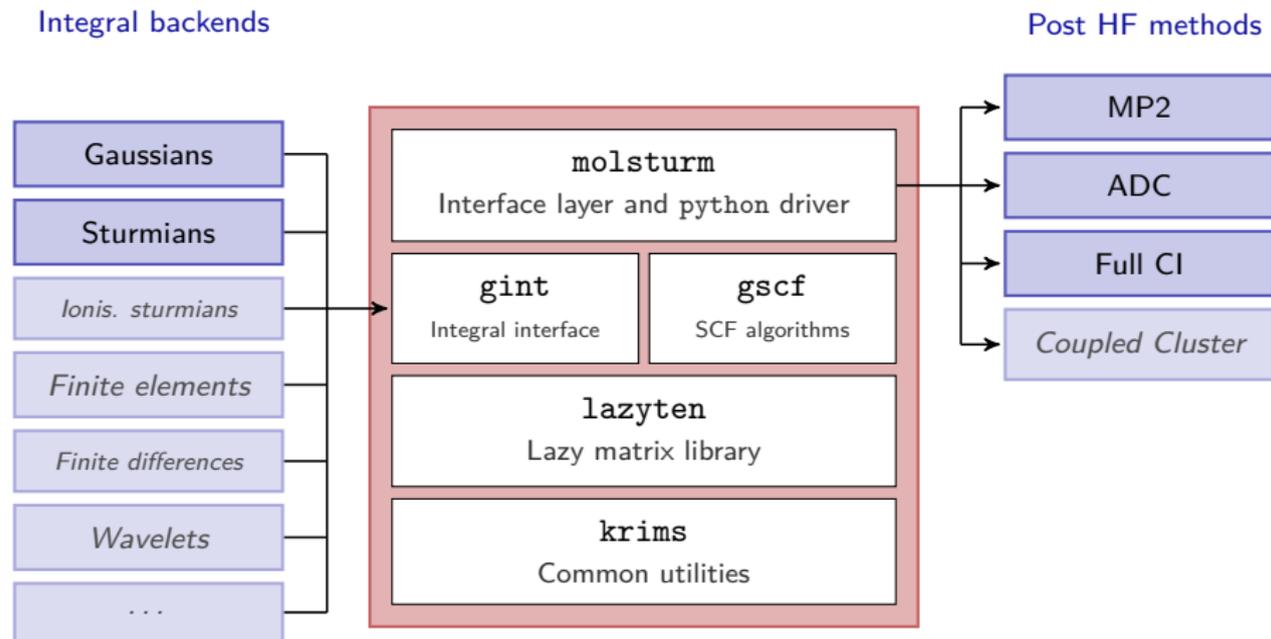
⇒ Explore methods across basis function types<sup>1,2</sup>

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<sup>1</sup>M. F. Herbst, A. Dreuw and J. E. Avery. J. Chem. Phys., **149**, 84106 (2018)

<sup>2</sup>M. F. Herbst. Ph.D. thesis, Ruprecht-Karls-Universität Heidelberg (2018)

# molsturm structure



## Principle ingredient: Contraction-based methods

- Contraction-based methods
  - Avoid **storing** matrices
  - Employ iterative, subspace-based algorithms
  - **Contraction** expressions (e.g. matrix-vector products)
  - Common in Post-HF: *Working equations*

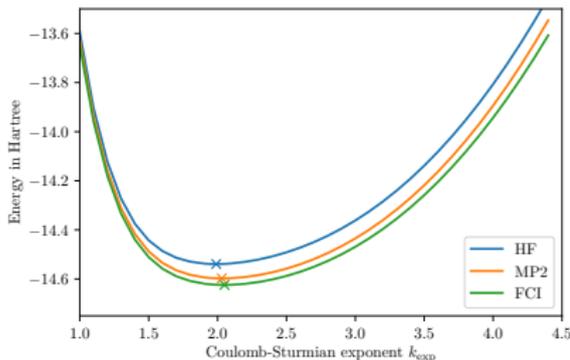
⇒ SCF code only needs Fock contraction

⇒ Hide discretisation details inside Fock object

⇒ Flexible to exploit discretisation-specific properties

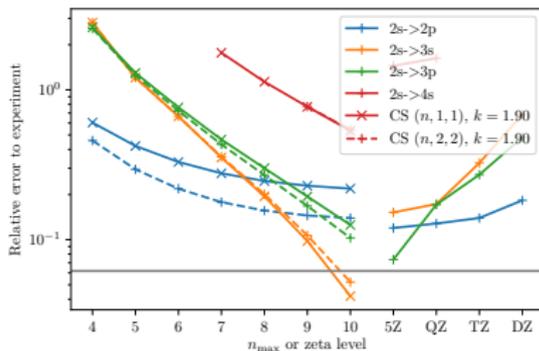
## molsturm interface: Linked codes

Coulomb-Sturmian based MP2 and FCI



- FCI from `pyscf`<sup>1</sup>
- Coulomb Sturmians from `sturmint`<sup>2</sup>

Coulomb-Sturmian and Gaussian based ADC(2)



- ADC(2) from `adcmans`<sup>3</sup>
- Gaussians from `libint`<sup>4</sup> or `libcint`<sup>5</sup>

<sup>1</sup>Q. Sun et al. WIREs Comput Mol Sci, **8**, e1340 (2017).

<sup>2</sup>J. E. Avery and M. F. Herbst. <https://molsturm.org/sturmint> (2018)

<sup>3</sup>M. Wormit et al. Mol. Phys., **112**, 774 (2014).

<sup>4</sup>E. Valeev et al. *evaleev/libint: 2.3.1* (2017).

<sup>5</sup>Q. Sun. J. Comput. Chem., **36**, 1664 (2015)