

NEW APPROACHES TOWARDS A FINITE-ELEMENT BASED HARTREE-FOCK METHOD

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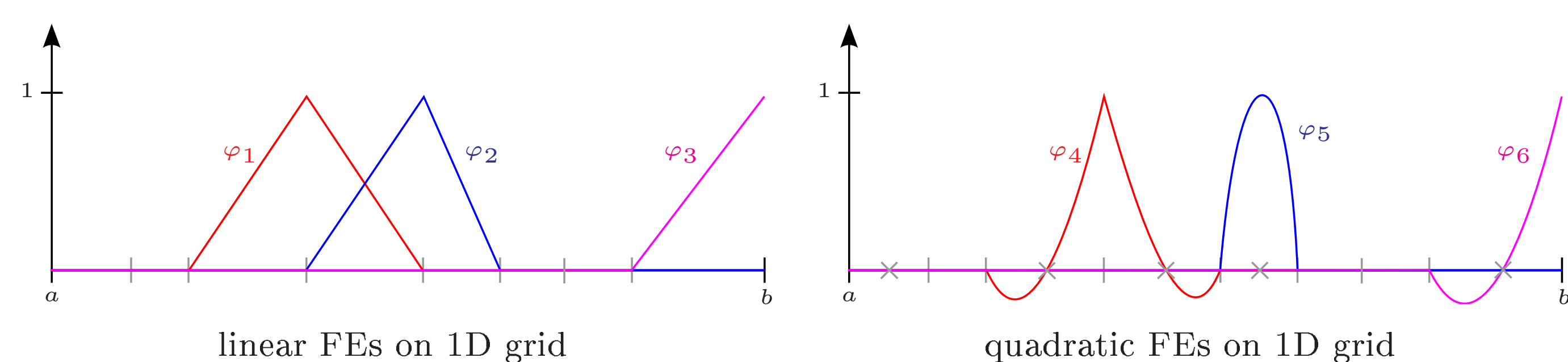
INTRODUCTION

Employing a finite-element basis in electronic structure calculations provides a novel alternative to the commonly used atom-centered (AC) basis functions [1, 2]. Expected properties:

- Sparse, nearly diagonal-dominant matrices
- Easy parallelisation of code
- Adaptive refinement: Grid adapts to molecular structure automatically
- Intrinsic multiscale methods possible
- Reduction of implicit bias imposed by AC basis
- More basis functions required compared to a pure AC basis

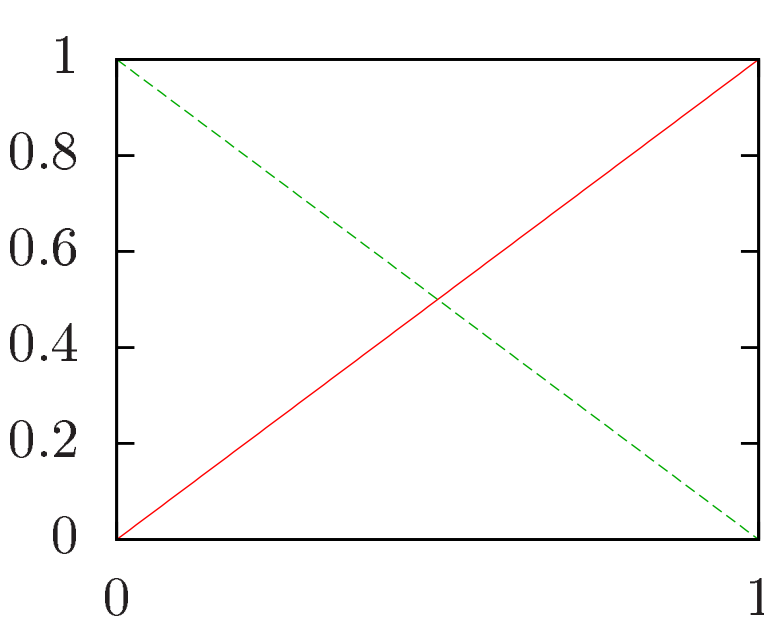
FINITE ELEMENTS (FES)

- Piecewise polynomials $\{\varphi_j\}$ on discretised grid
- Only non-zero in a few grid cells \Rightarrow strongly **localised**
- At **nodal points** $\{\mathbf{x}_k\}$: $\varphi_j(\mathbf{x}_k) = \delta_{jk}$
- Hence: $f(x) = \sum_j z_j \varphi_j \Rightarrow f(x_k) = z_k$

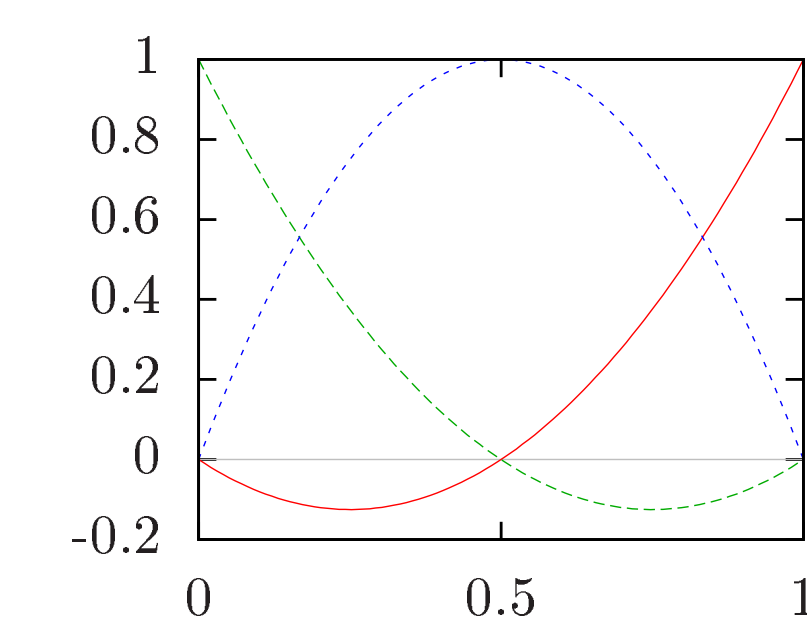


- **Affine map** $\mu_c : c_0 \rightarrow c$, between arbitrary cells on discretisation grid c and reference cell c_0
- Allows description of FEs by only considering a small number of **shape functions** $\{e_\alpha\}$ (Lagrange polynomials), defined on c_0 :

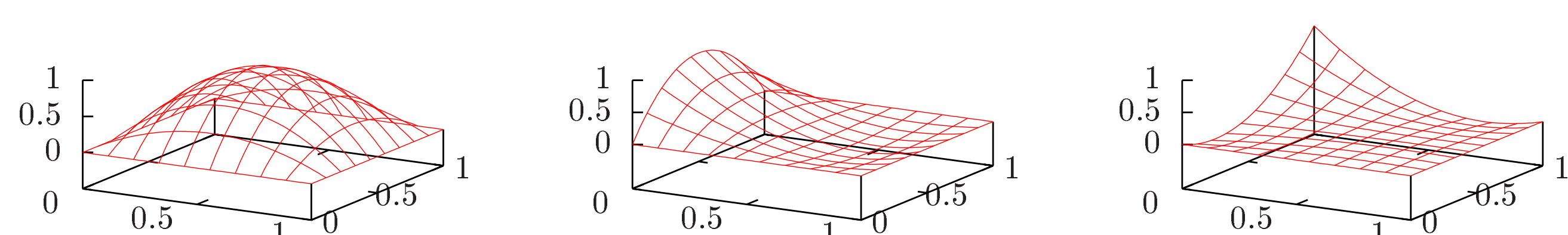
$$\varphi_k|_c(\mathbf{r}) = e_\alpha(\mu_c^{-1}(\mathbf{r})) \quad \text{for some } \alpha$$



1D linear shape functions



1D quadratic shape functions



2D quadratic shape functions

- **A posteriori error estimation**: Estimate cell-wise error bound after calculation
- **Adaptive grid refinement**: Construct new grid, where cells with largest *a posteriori* error are refined
- **Intrinsic multiscale**: Refine grid in particular regions of space which are of special interest

WEAK FORMULATION OF HARTREE-FOCK

- Consider orbital in FE basis: $\psi_i = \sum_j z_j^{(i)} \varphi_j$
- Define:

$$a(f, g) := \int_{\Omega} \frac{1}{2} \nabla f(\mathbf{r}) \cdot \nabla g(\mathbf{r}) + f(\mathbf{r}) V(\mathbf{r}) g(\mathbf{r}) \, d\mathbf{r} \quad (1)$$

$$V = \underbrace{V_0}_{\text{electron-nuclear interaction}} + \underbrace{V_H}_{\text{Hartree potential}} + \underbrace{V_x}_{\text{exchange potential}} \quad (2)$$

$$m(f, g) := \int_{\Omega} f(\mathbf{r}) \cdot g(\mathbf{r}) \, d\mathbf{r} \quad (3)$$

- We solve the Hartree-Fock (HF) equation in its **weak form**

$$\begin{aligned} \forall \varphi_j : a(\varphi_j, \psi_i) &= \varepsilon_i m(\varphi_j, \psi_i) \\ \Rightarrow \mathbf{A} \mathbf{z}^{(i)} &= \varepsilon_i \mathbf{M} \mathbf{z}^{(i)} \end{aligned} \quad (4)$$

where $A_{jk} = a(\varphi_j, \varphi_k)$ and $M_{jk} = m(\varphi_j, \varphi_k)$.

- Solving (4) in an FE basis is usually difficult due to the high dimensionality

METHOD OVERVIEW

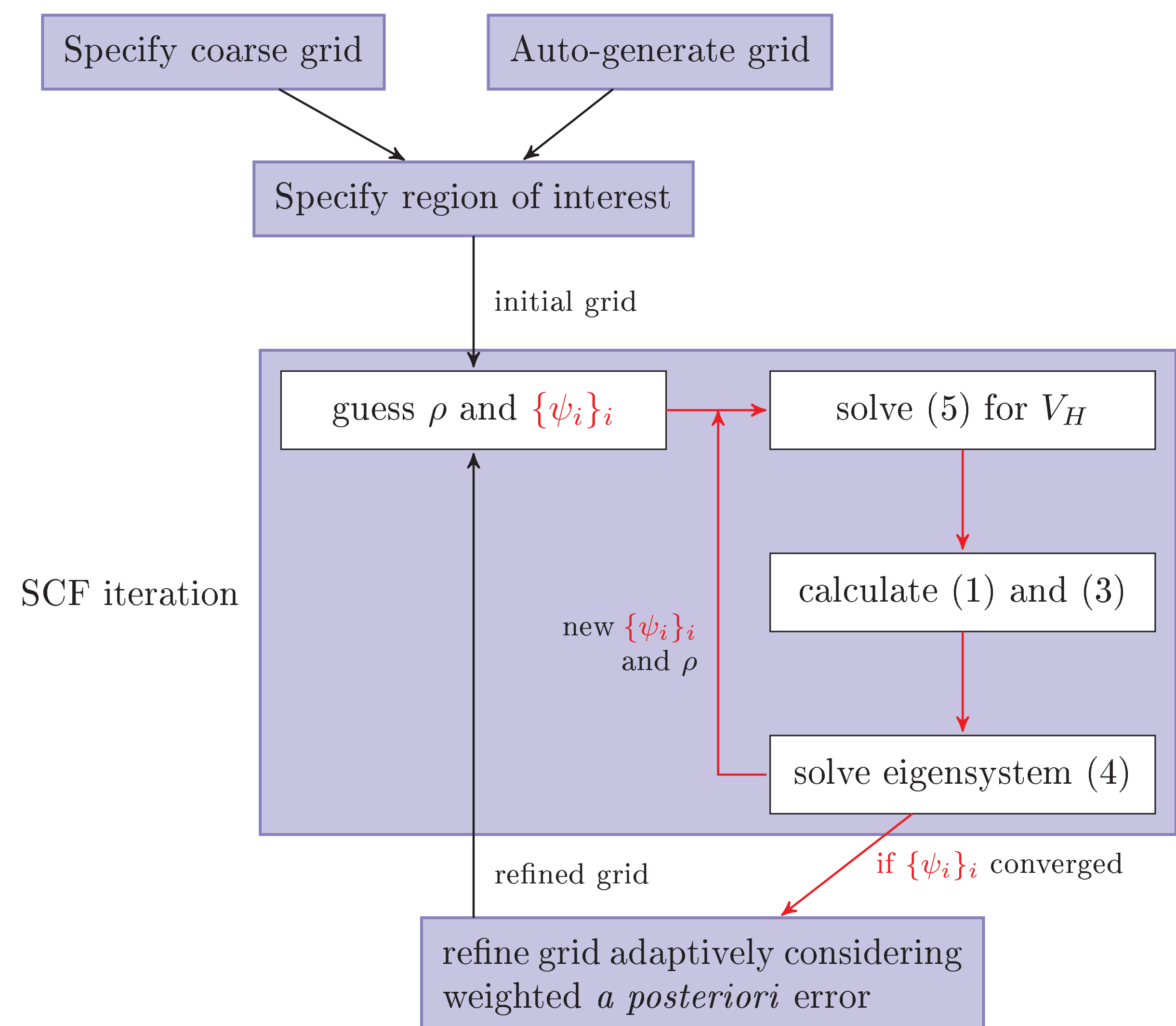
- Based on **SIESTA** method [3, 4]
- FE basis allows calculation of Hartree potential V_H by Poisson equation

$$-\Delta V_H(\mathbf{r}) = \rho(\mathbf{r}) \quad (5)$$

in **linear time**.

- To avoid dimensionality issues in the solution of (4) represent core orbitals in an AC basis
- Still treat valence orbitals in non-biased FE basis \Rightarrow **reduction of bias** by AC basis
- Fit AC basis functions to FE basis functions such that core orbitals are available in both representations
- Non-local Hartree-Fock exchange potential problematic
- Back end for FE: **deal.ii library** [5]

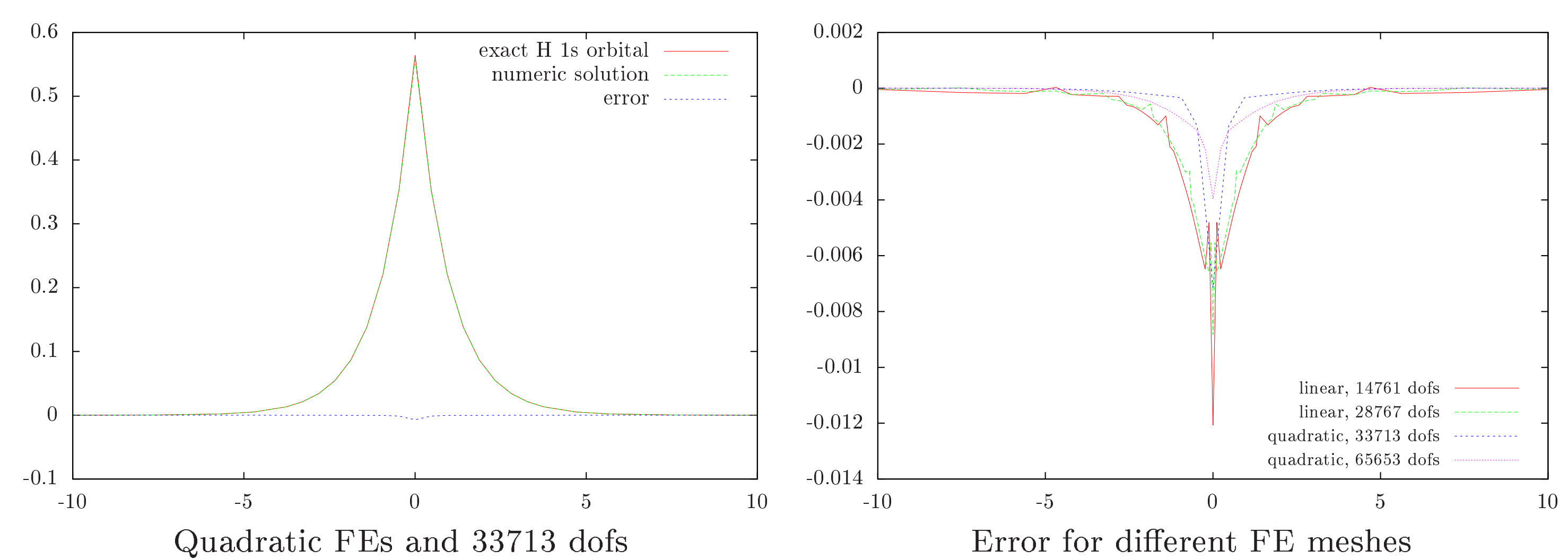
PROPOSED CALCULATION SCHEME



- Can also skip grid refinement and re-use grid from previous calculation

SOLUTION FOR HYDROGEN: 1S ORBITAL

- Comparison of analytic solution with solutions on different FE meshes plotted along the x axis



- Greatest error at the discontinuity of ψ_i near the nucleus

OUTLOOK

- Use ECPs to get rid of discontinuity at nucleus
- Develop an approximate Fock operator in the FE context to avoid using exact one each SCF iteration
- Implement automatic mesh generation and an intrinsic multiscale method

REFERENCES

- [1] J. E. Pask and P. A. Sterne, *Modell. Simul. Mater. Sci. Eng.*, **13** (2005), R71.
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